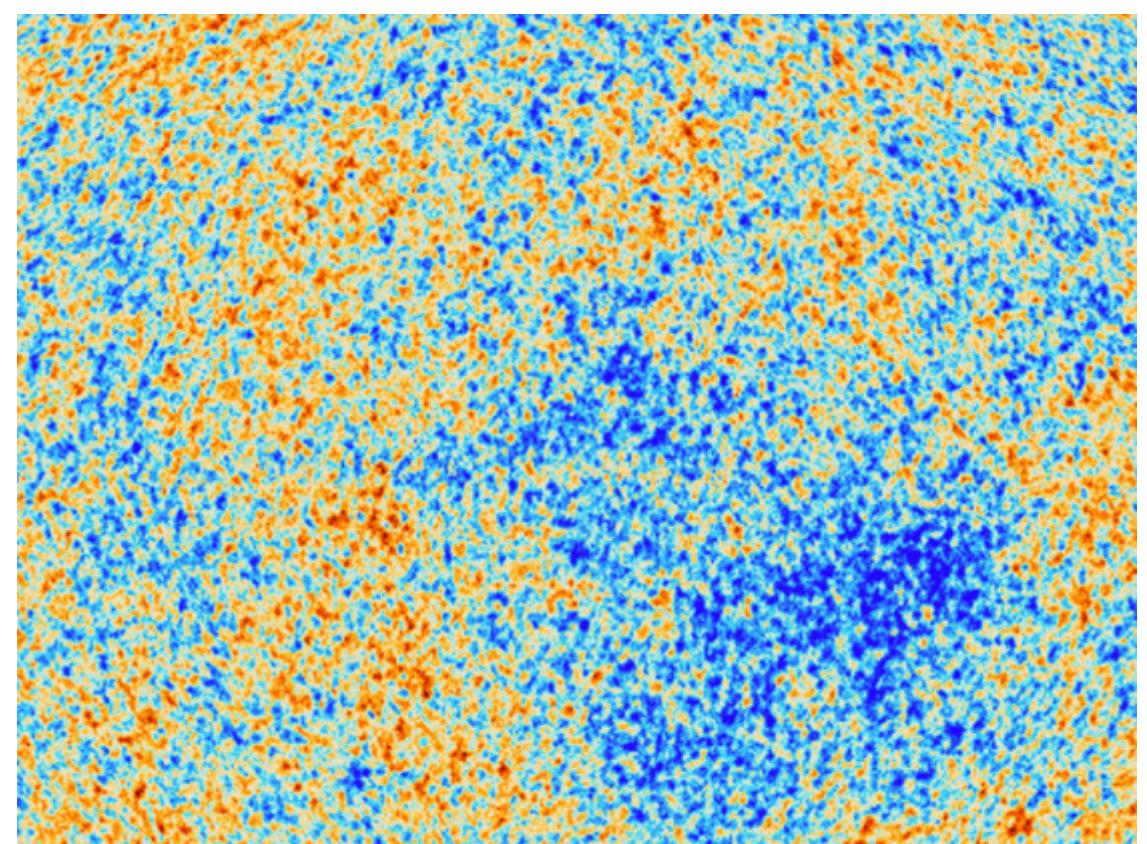
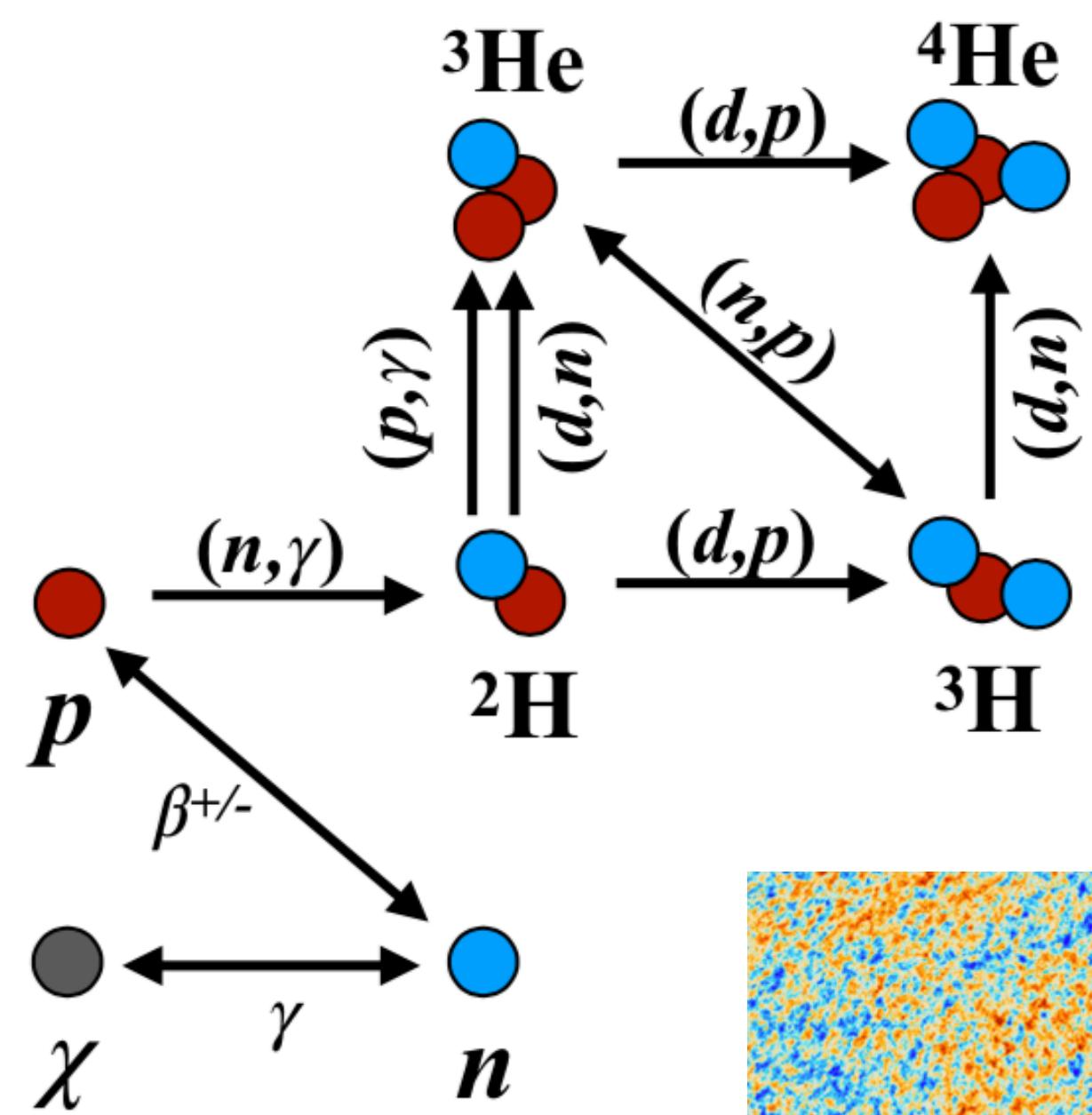


Dark Neutrons: Cosmology & Astrophysics

Nirmal Raj
TRIUMF

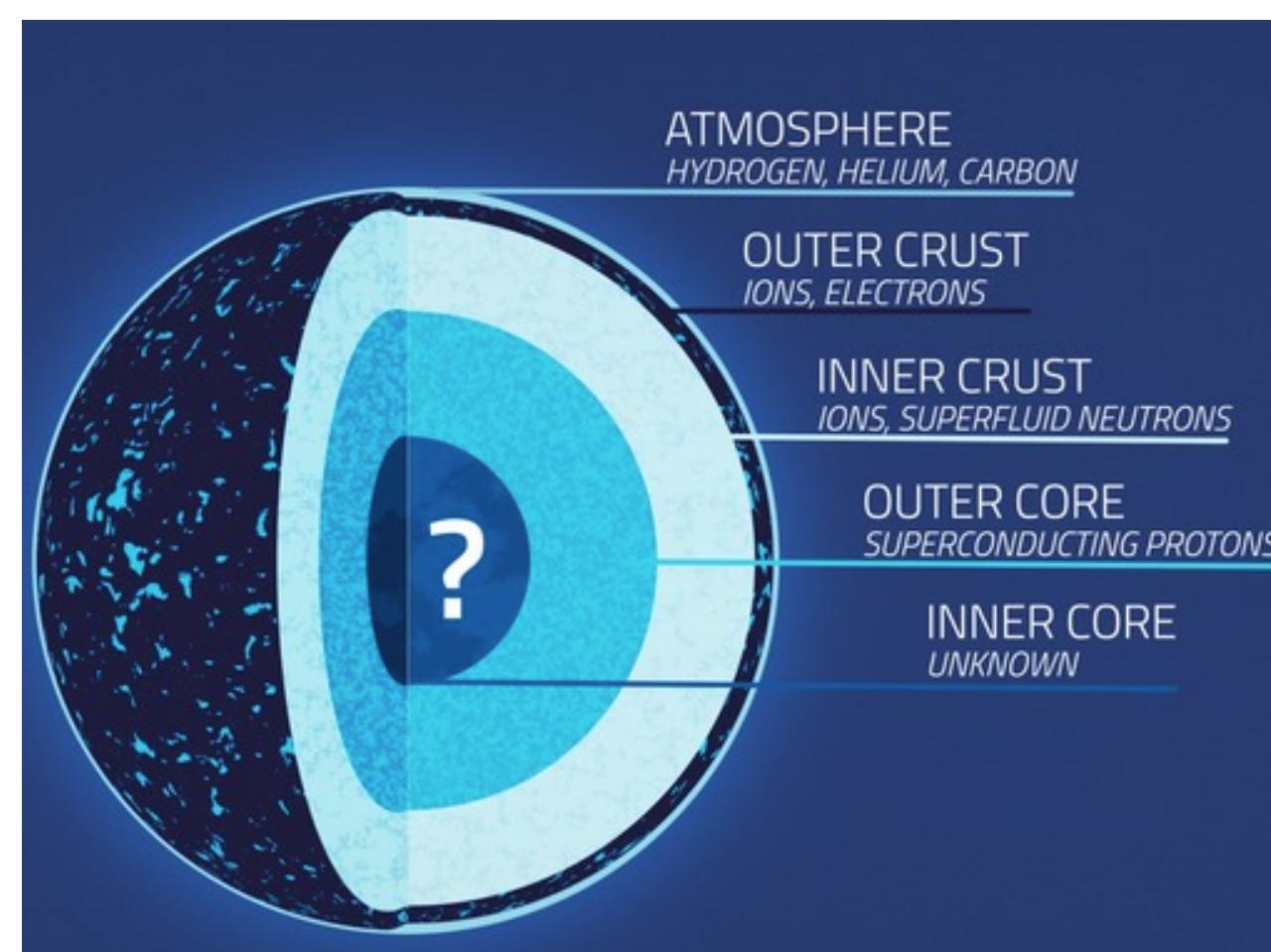
Phys. Rev. Lett. 125 (2020), 231803,
Phys. Rev. Lett. 127 (2021), 061805,
Phys. Rev. D. (2021) 103.115002

with David McKeen
& Maxim Pospelov



Brookhaven
Forum: *Opening
New Windows to
the Universe*

Nov 03, 2021



Introduction

hypothesis: a new particle “ χ ”

its character: 0 : charge under all fundamental forces

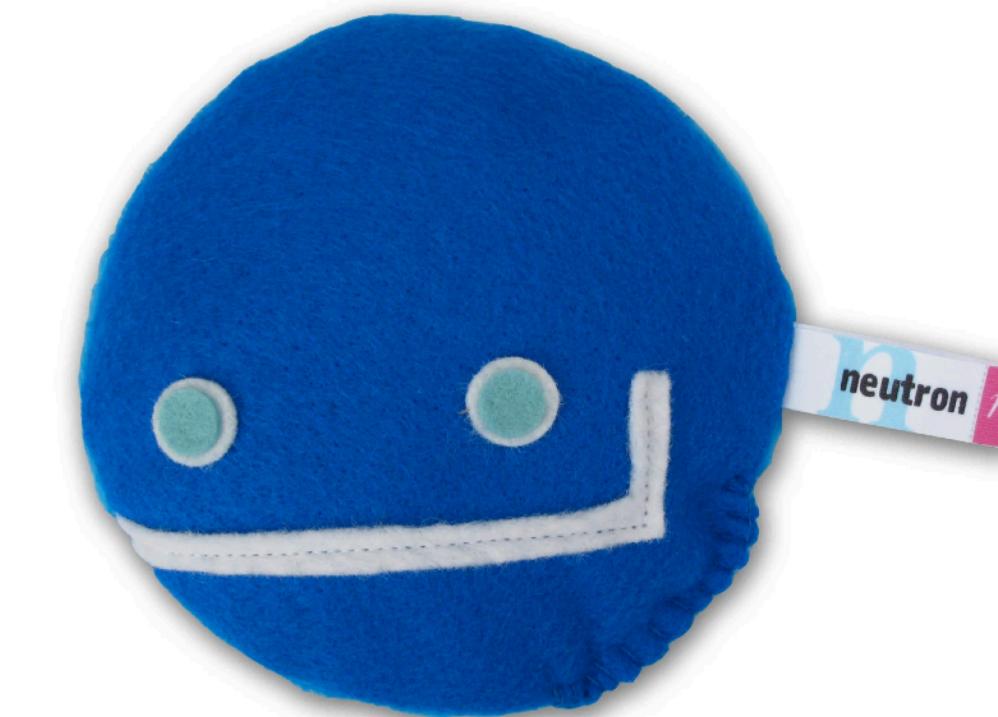
1/2 : spin

1 : baryon number



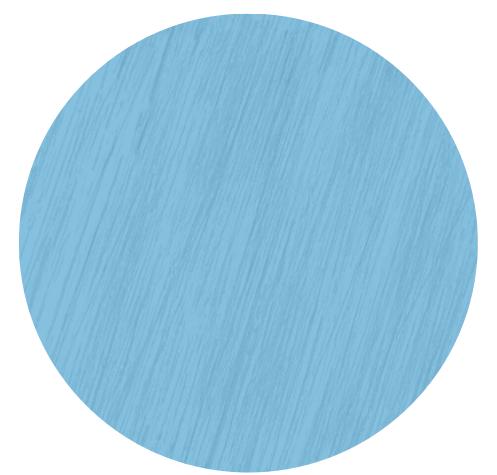
James Chadwick

It's called a neutron.
N. E. U. T. R. O. N,
neutron.



also $\Lambda^0, \Sigma^0, \Delta^0, \dots$

neutron



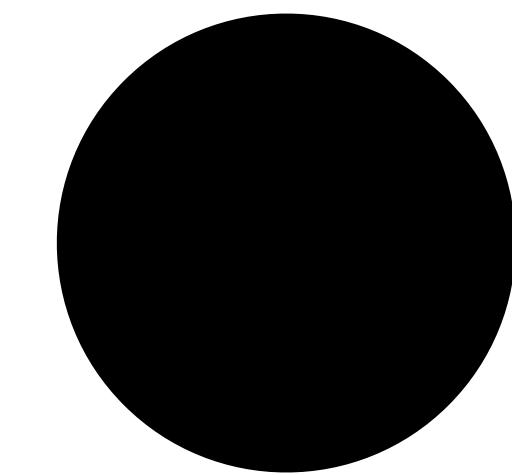
n

m_n



939.5654 MeV/ c^2

“dark” neutron



χ

m_χ



?

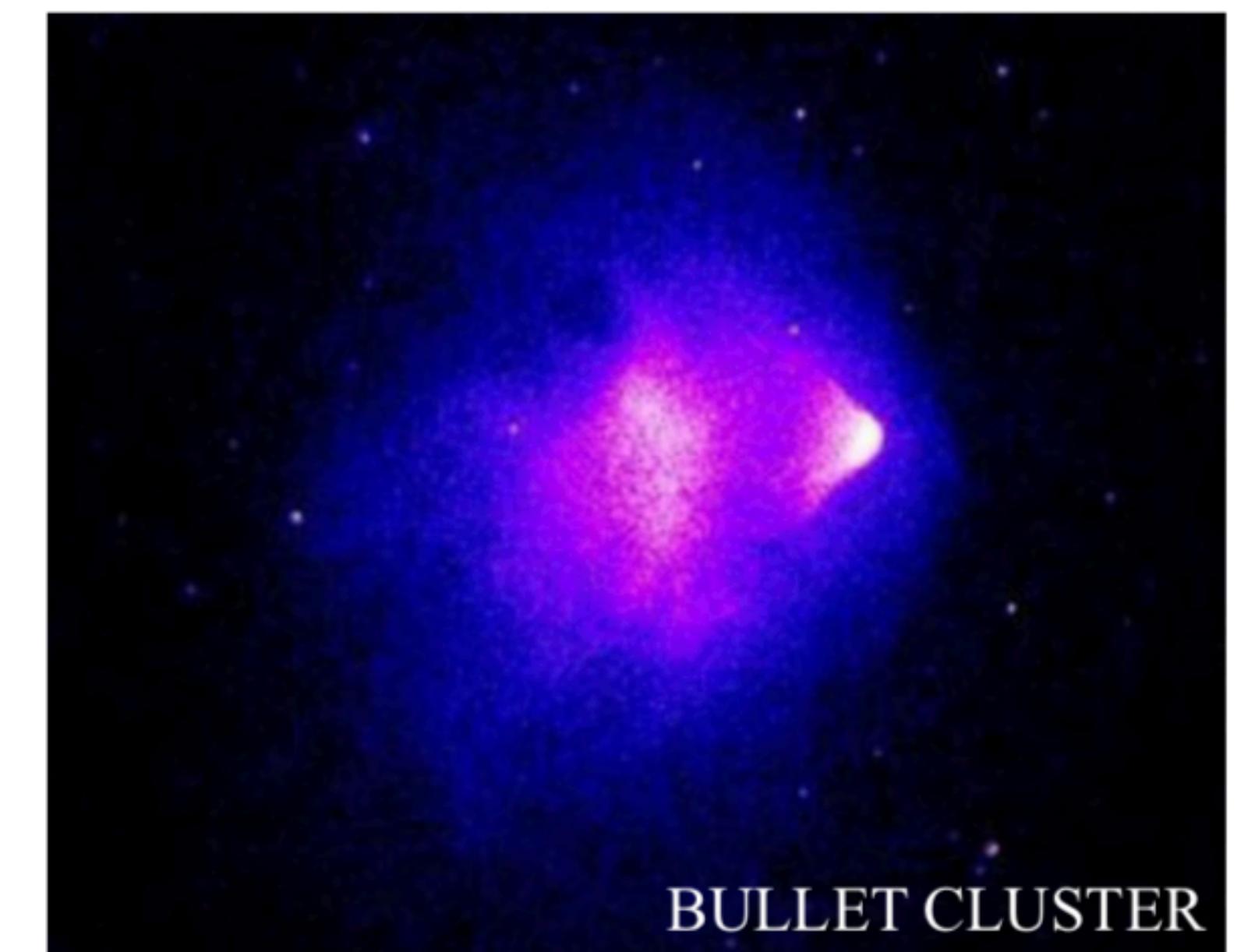
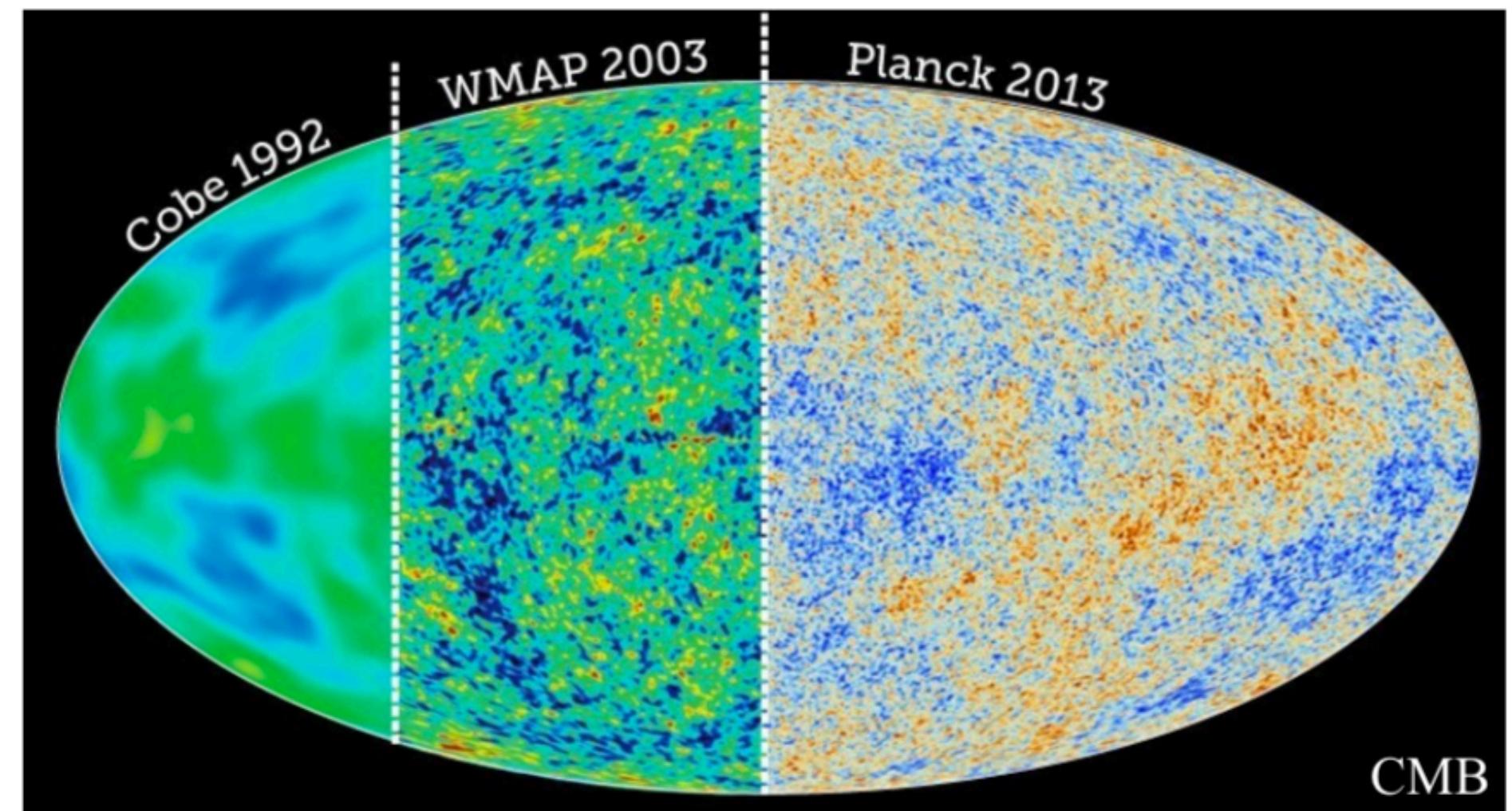
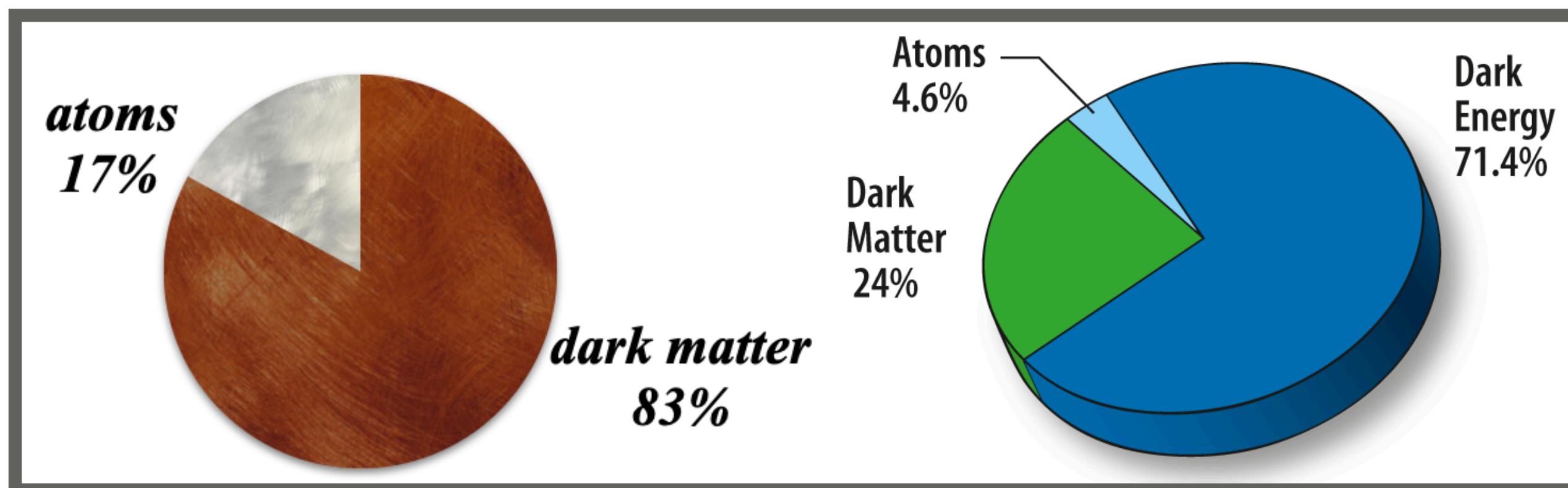
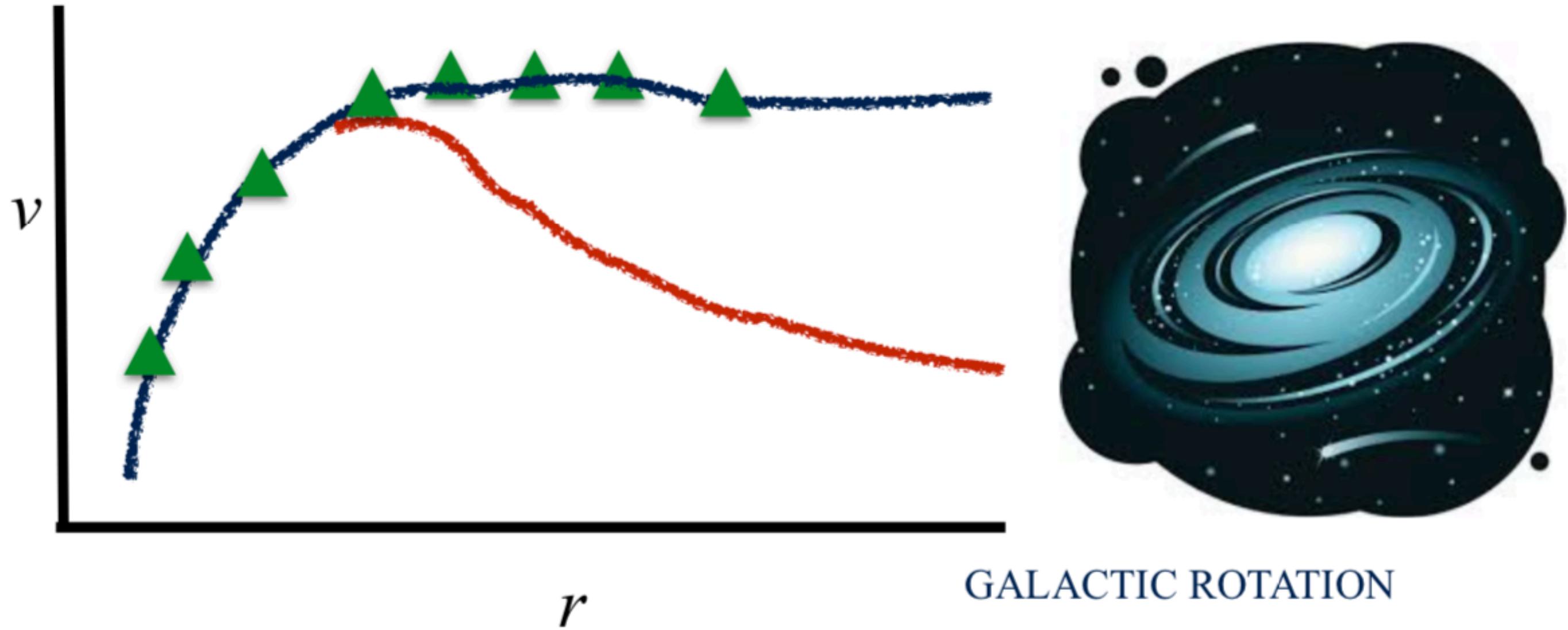
Hamiltonian

$$\begin{pmatrix} \bar{m}_n & \epsilon_{n\chi} \\ \epsilon_{n\chi} & \bar{m}_\chi \end{pmatrix}$$

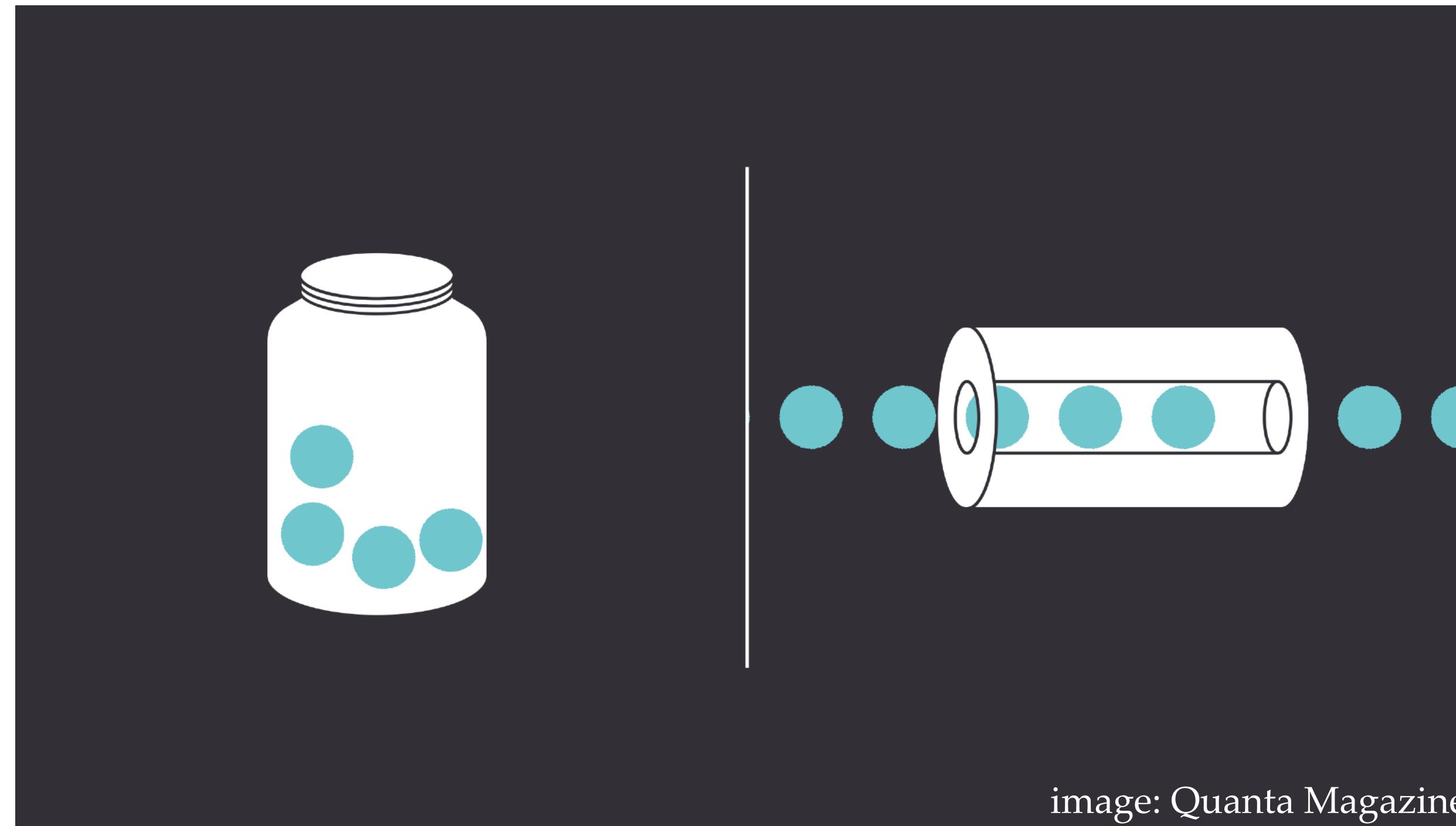
The Hamiltonian matrix is shown in parentheses. It has two columns and two rows. The top-left entry is \bar{m}_n , the top-right is $\epsilon_{n\chi}$, the bottom-left is $\epsilon_{n\chi}$, and the bottom-right is \bar{m}_χ . Above the matrix, there are two sets of three dots: one set above the first column and one set above the second column, indicating continuation.

Why care?

(1) the *dark matter* of the universe



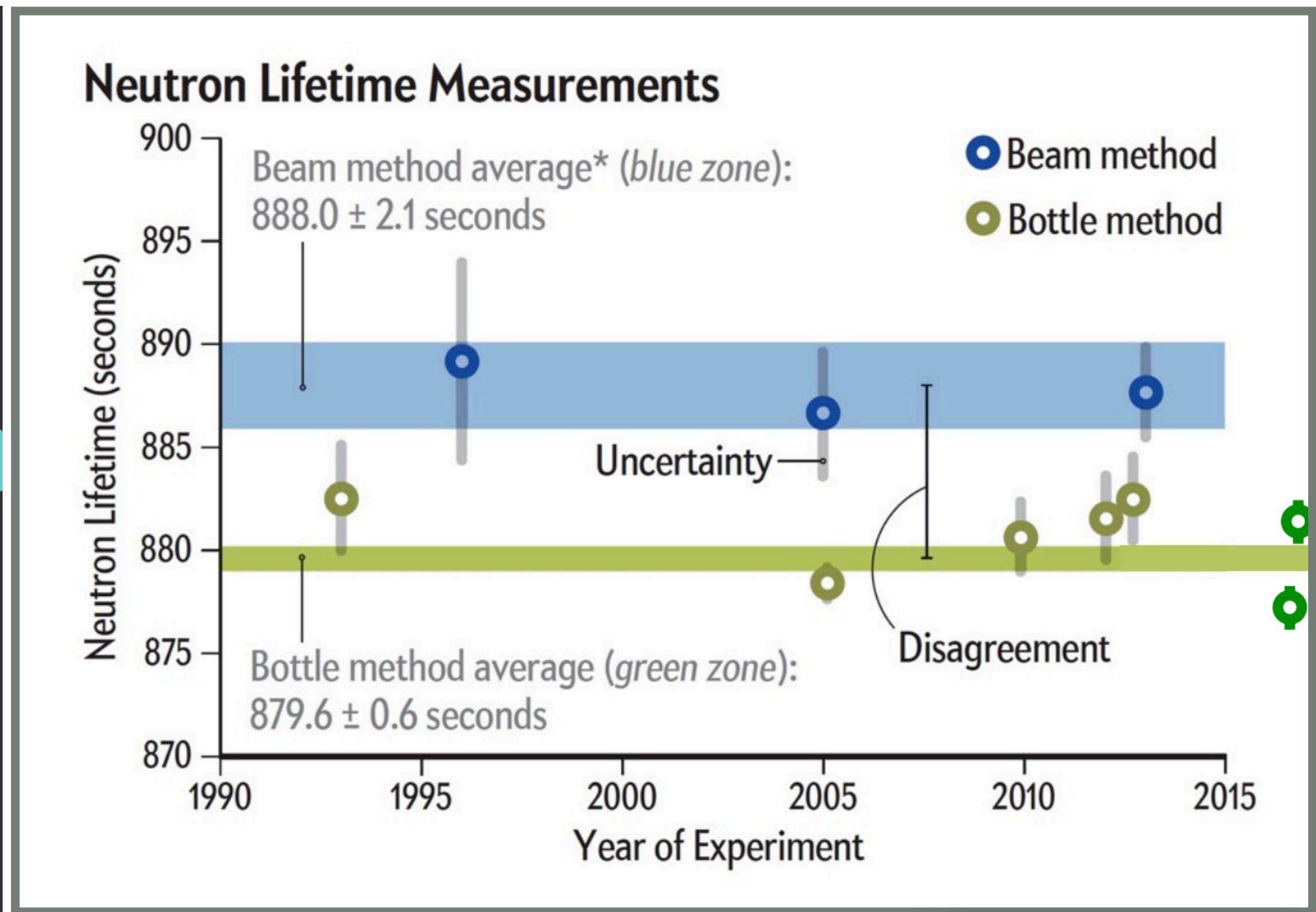
Why care? (2) the *neutron lifetime puzzle*



explain puzzle with

1% branching to
 $n \rightarrow \chi + \text{anything in bottle}$
Fornal, Grinstein (2018)

1% probability of
 $n \rightarrow \chi$ in **beam**
Berezhiani (2018)



discrepancy: $\frac{\Delta\tau_n}{\tau_n} \approx 1\%$

Why care?

(3) the “*XENON1T excess*” from last summer

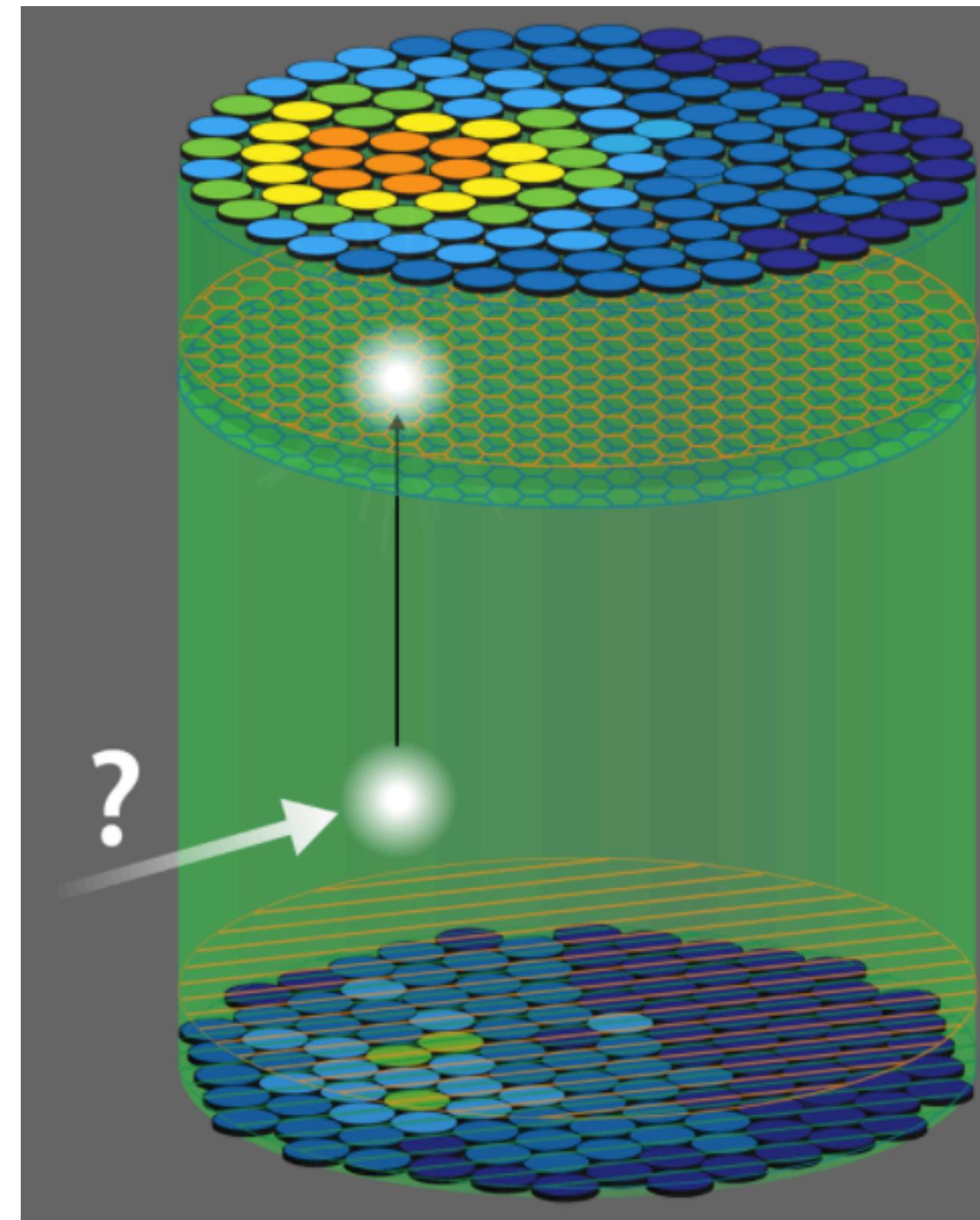
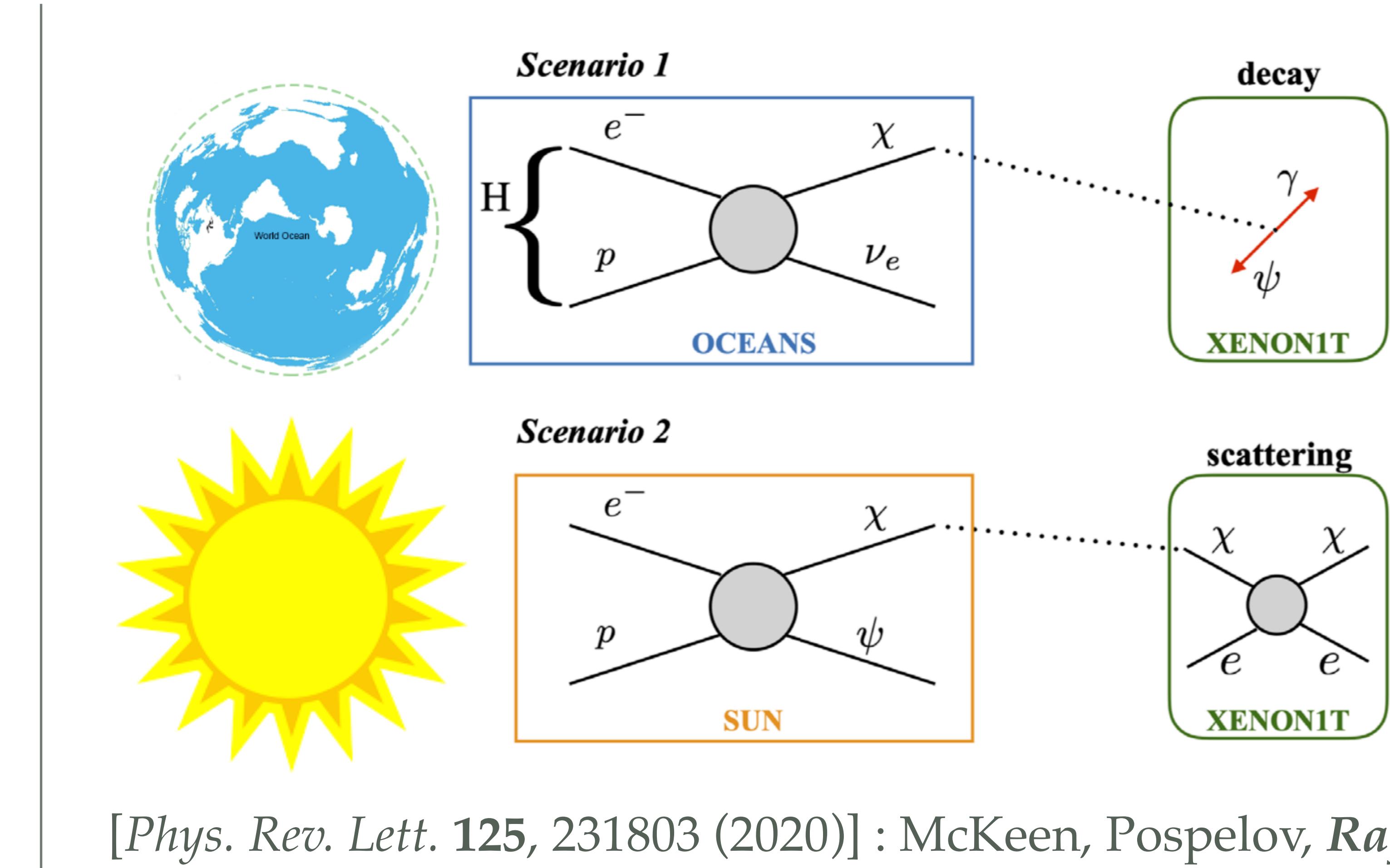


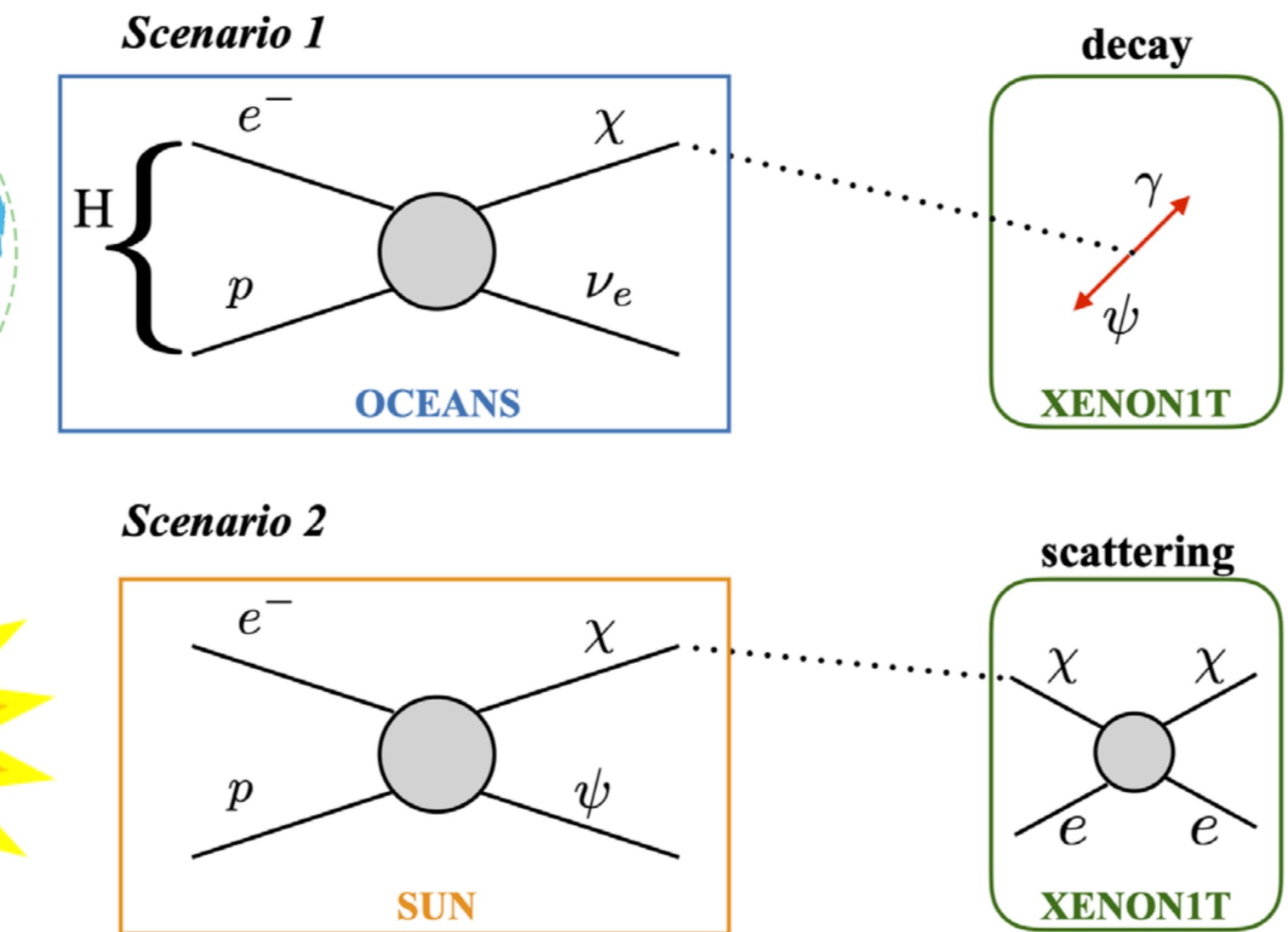
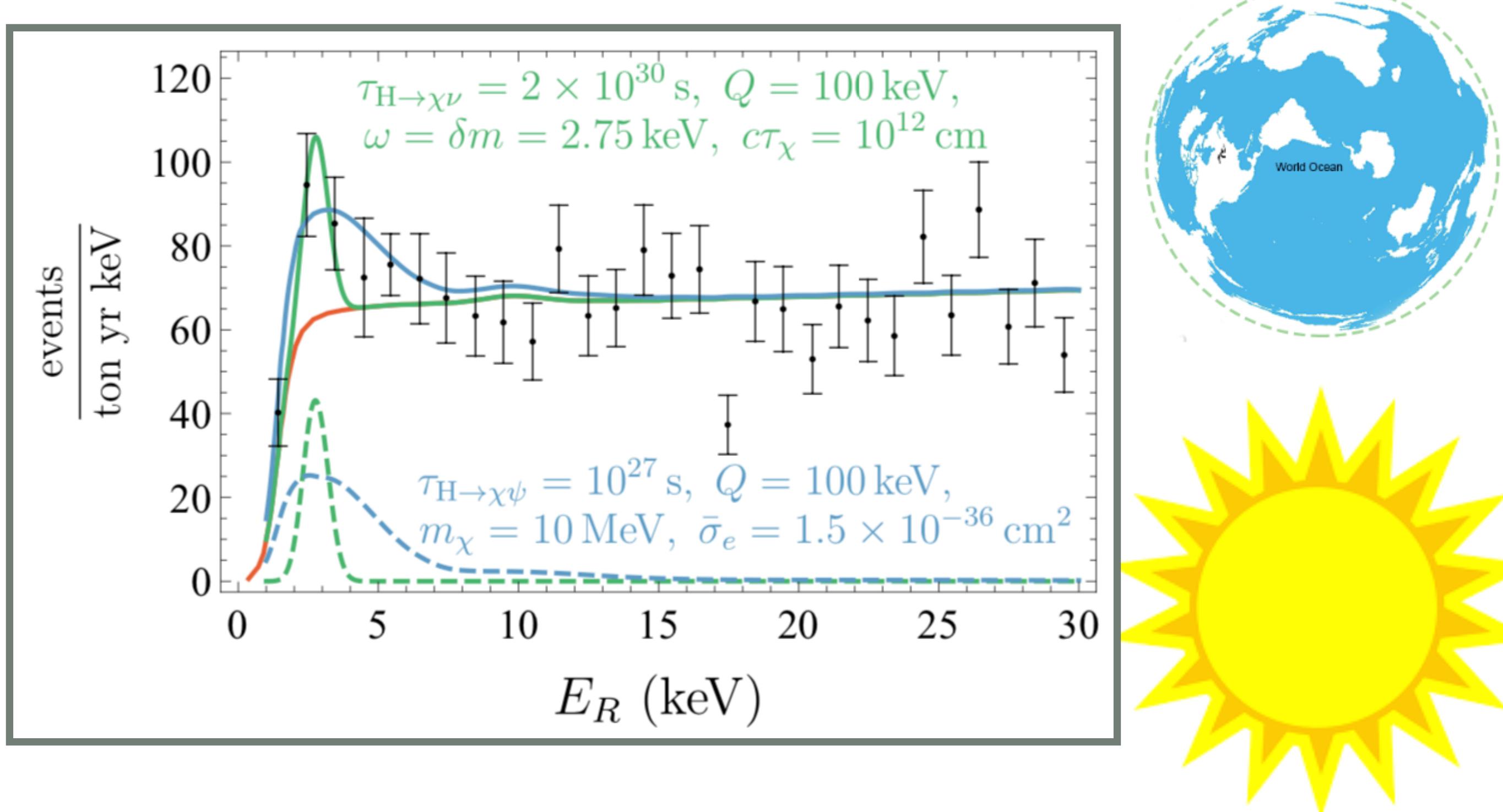
image: APS

arXiv: 2006.09721



Why care?

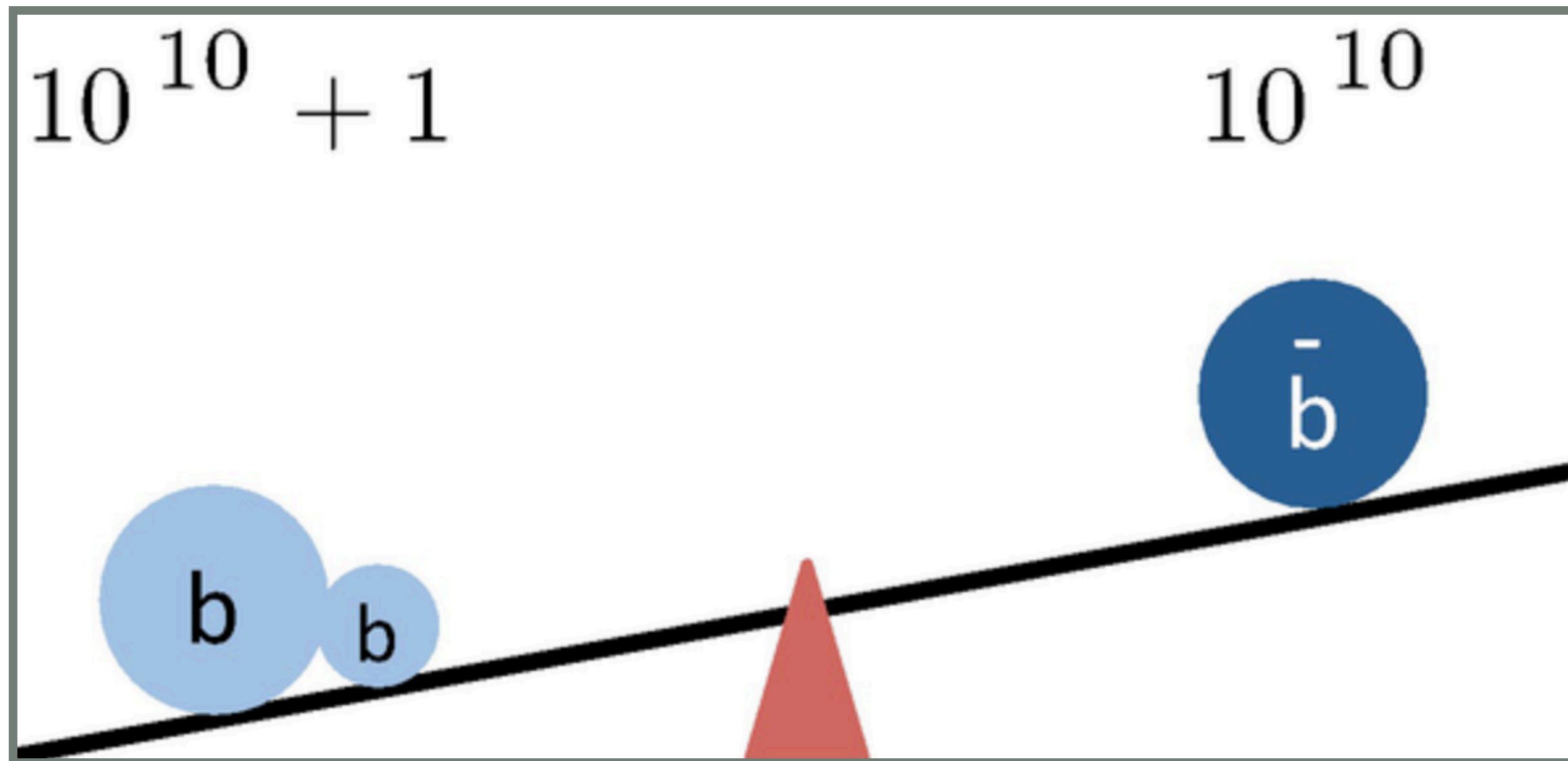
(3) the “*XENON1T excess*” from last summer



[*Phys. Rev. Lett.* **125**, 231803 (2020)] : McKeen, Pospelov, Raj

Why care?

(4) role in baryon asymmetry



D. McKeen and A. E. Nelson, Phys. Rev. D **94**, 076002 (2016), arXiv:1512.05359 [hep-ph].

K. Aitken, D. McKeen, T. Neder, and A. E. Nelson, Phys. Rev. D **96**, 075009 (2017), arXiv:1708.01259 [hep-ph].

K. Babu, P. Bhupal Dev, E. C. Fortes, and R. Mohapatra, Phys. Rev. D **87**, 115019 (2013), arXiv:1303.6918 [hep-ph]; R. Allahverdi, P. S. B. Dev, and B. Dutta, Phys. Lett. B **779**, 262 (2018), arXiv:1712.02713 [hep-ph]; G. Elor, M. Escudero, and A. Nelson, Phys. Rev. D **99**, 035031 (2019), arXiv:1810.00880 [hep-ph]; A. E. Nelson and H. Xiao, Phys. Rev. D **100**, 075002 (2019), arXiv:1901.08141 [hep-ph]; G. Alonso-Álvarez, G. Elor, A. E. Nelson, and H. Xiao, JHEP **03**, 046 (2020), arXiv:1907.10612 [hep-ph].

T. Bringmann, J. M. Cline, and J. M. Cornell, Phys. Rev. D **99**, 035024 (2019), arXiv:1810.08215 [hep-ph].

Model

@ hadron level :

$$\mathcal{L} \supset -\delta(\bar{\chi}n + \bar{n}\chi)$$

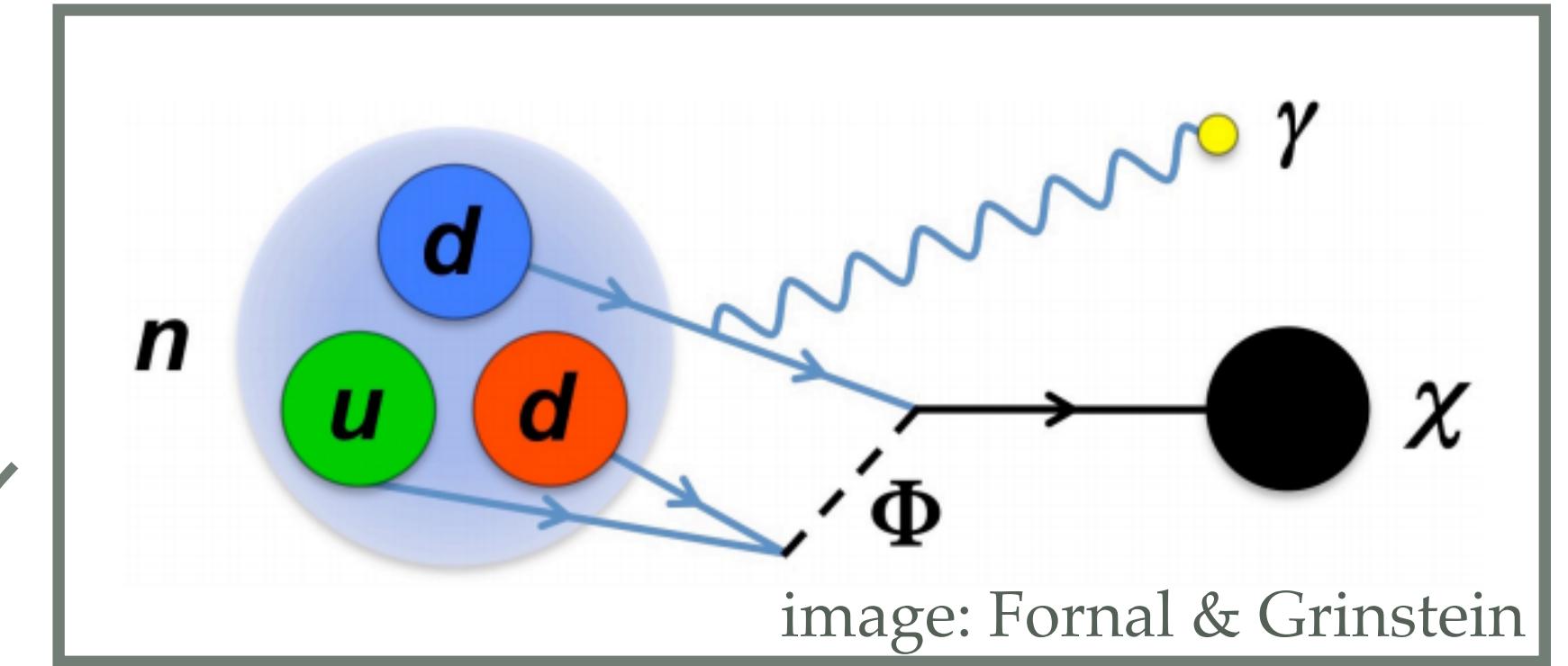
$$\mathcal{L}_{\text{eff}} \supset \frac{\mu_n}{2} \theta \bar{\chi} \sigma^{\mu\nu} n F_{\mu\nu} + \text{h.c.}$$

$$\mu_n = -1.91 \mu_N$$

neutron magnetic moment

$$\left(\begin{array}{l} \mu_N = e/(2m_n) \simeq 0.1 \text{ e fm} \\ \text{nuclear magneton} \end{array} \right)$$

exotic neutron decay



n lifetime puzzle:

$$\text{Br}_{n \rightarrow \chi\gamma} \simeq 0.01 \left(\frac{\theta}{5 \times 10^{-10}} \right)^2 \left(\frac{\Delta m}{\text{MeV}} \right)^3$$

$$\Gamma_{\chi \rightarrow n\gamma} \simeq \frac{1}{2200 \text{ s}} \left(\frac{\theta}{10^{-10}} \right)^2 \left| \frac{\Delta m}{10 \text{ MeV}} \right|^3$$

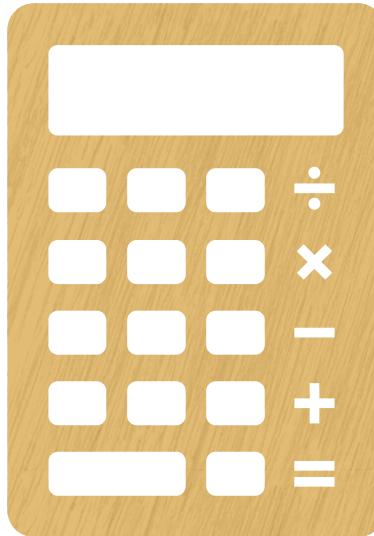
$$\Gamma_{\chi \rightarrow pe^- \bar{\nu}} = \frac{1}{9 \times 10^{22} \text{ s}} \left(\frac{\theta}{10^{-10}} \right)^2 \frac{F(Q_\chi/m_e)}{F(Q_n/m_e)}$$

Prehistoric census

Interesting cases:

$$(i) n_\chi^0 = 5.4(n_p^0 + n_n^0) \quad (\chi \text{ is the dark matter if } \tau_\chi > t_U)$$

$$(ii) n_\chi^0 = 0.01(n_p^0 + n_n^0) \quad (\text{perhaps never chem eqbm})$$



above QCD transition \Rightarrow quark level description required

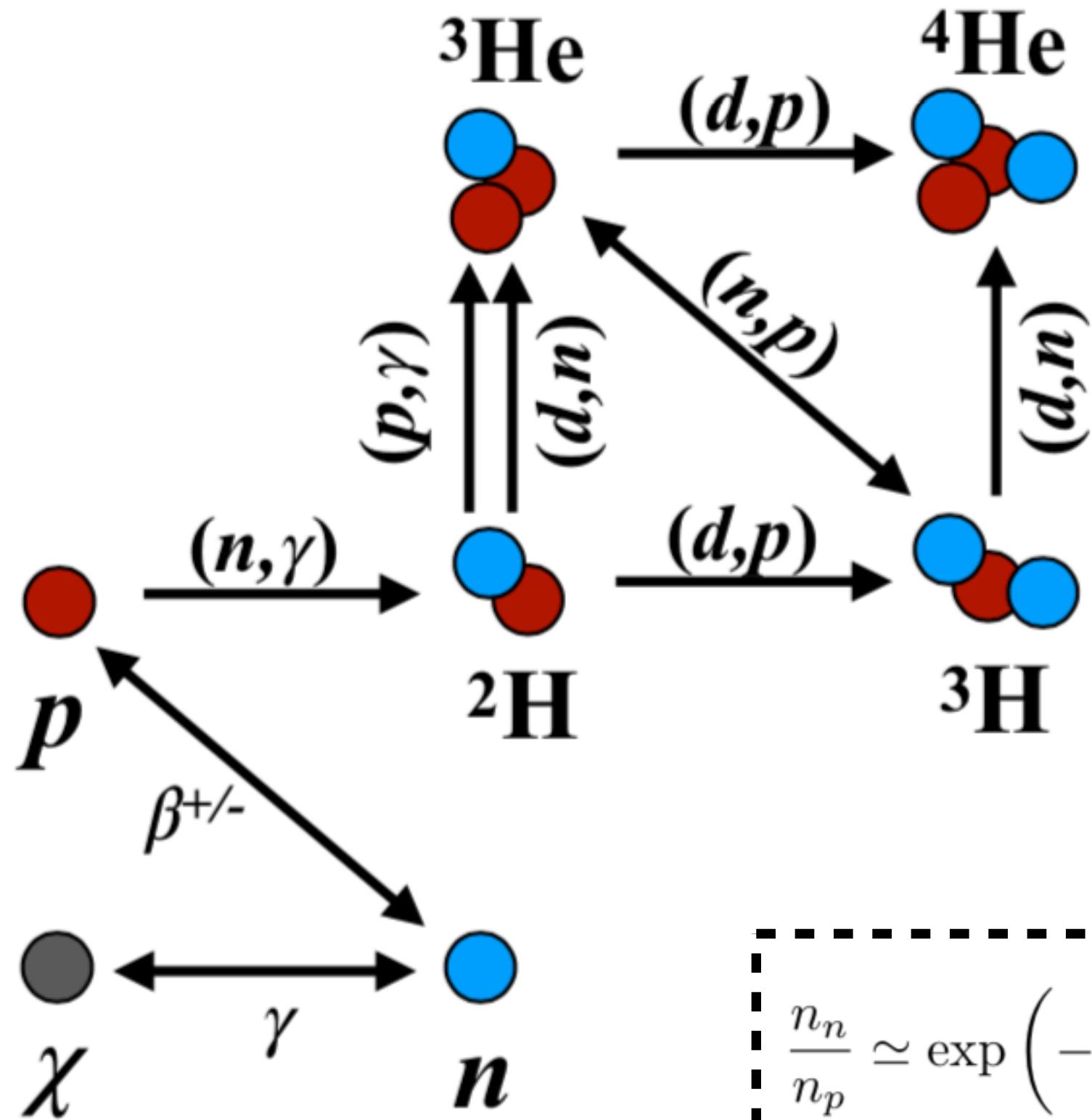
$$-\delta(\bar{\chi}n + \bar{n}\chi) \longleftrightarrow \bar{\chi}qqq/\Lambda^2 \Rightarrow \Gamma_{\Delta\chi} \sim T^5/\Lambda^4$$

chemical equilibrium keepable down to $T \sim \text{GeV—PeV}$

for $\theta \sim 10^{-20} — 10^{-10}$ and $\Delta m \sim 1 — 100 \text{ MeV}$

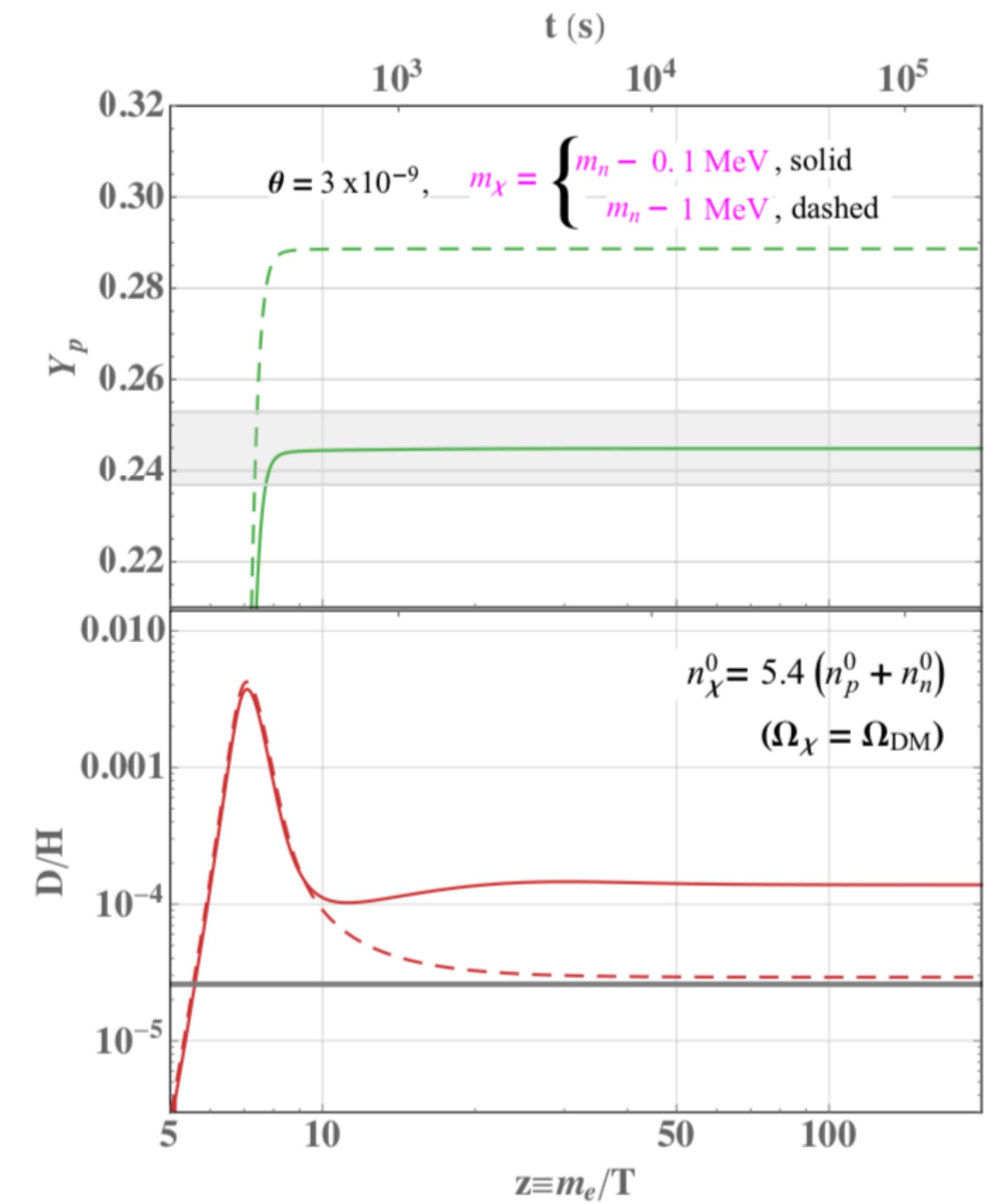
-----→ $n_\chi \sim n_p = n_n$ reasonable since universe was probably that hot

Probes: [1] primordial nucleosynthesis

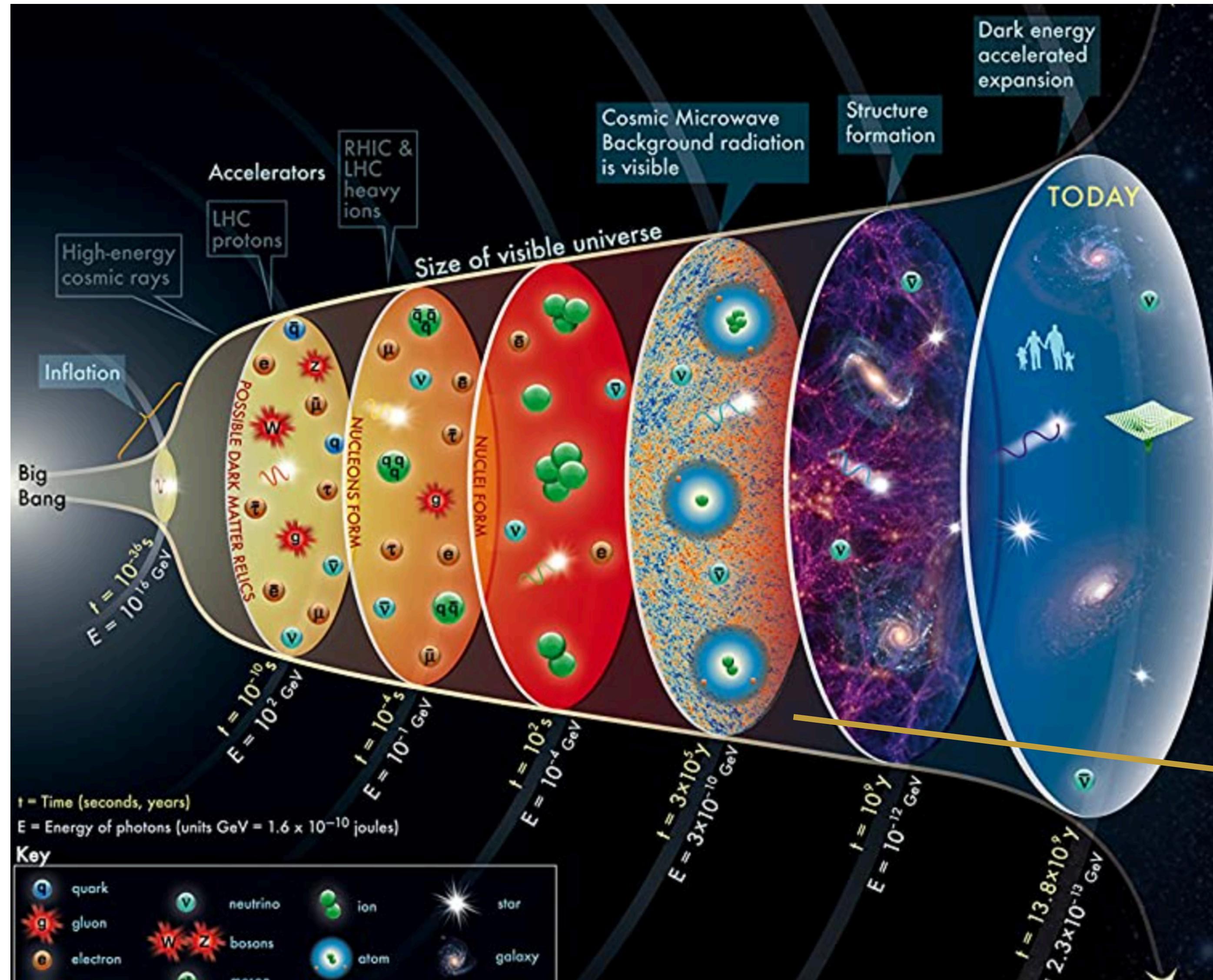


$$\frac{n_n}{n_p} \simeq \exp\left(-\frac{Q_{np}}{T_{np}}\right) \simeq \exp\left[-\frac{Q_{np}}{\bar{T}_{np}}\left(1 - \frac{\text{Br}_{n \rightarrow \chi}}{3}\right)\right]$$

$$\begin{aligned} \frac{\delta Y_p}{Y_p} &\simeq \frac{\delta(n_n/n_p)}{n_n/n_p} \times \frac{1}{1 + n_n/n_p} \\ &\simeq 0.4\% \left(\frac{\text{Br}_{n \rightarrow \chi}}{1\%}\right). \end{aligned}$$



Probes: [2] relic radiation



When kinematically open:

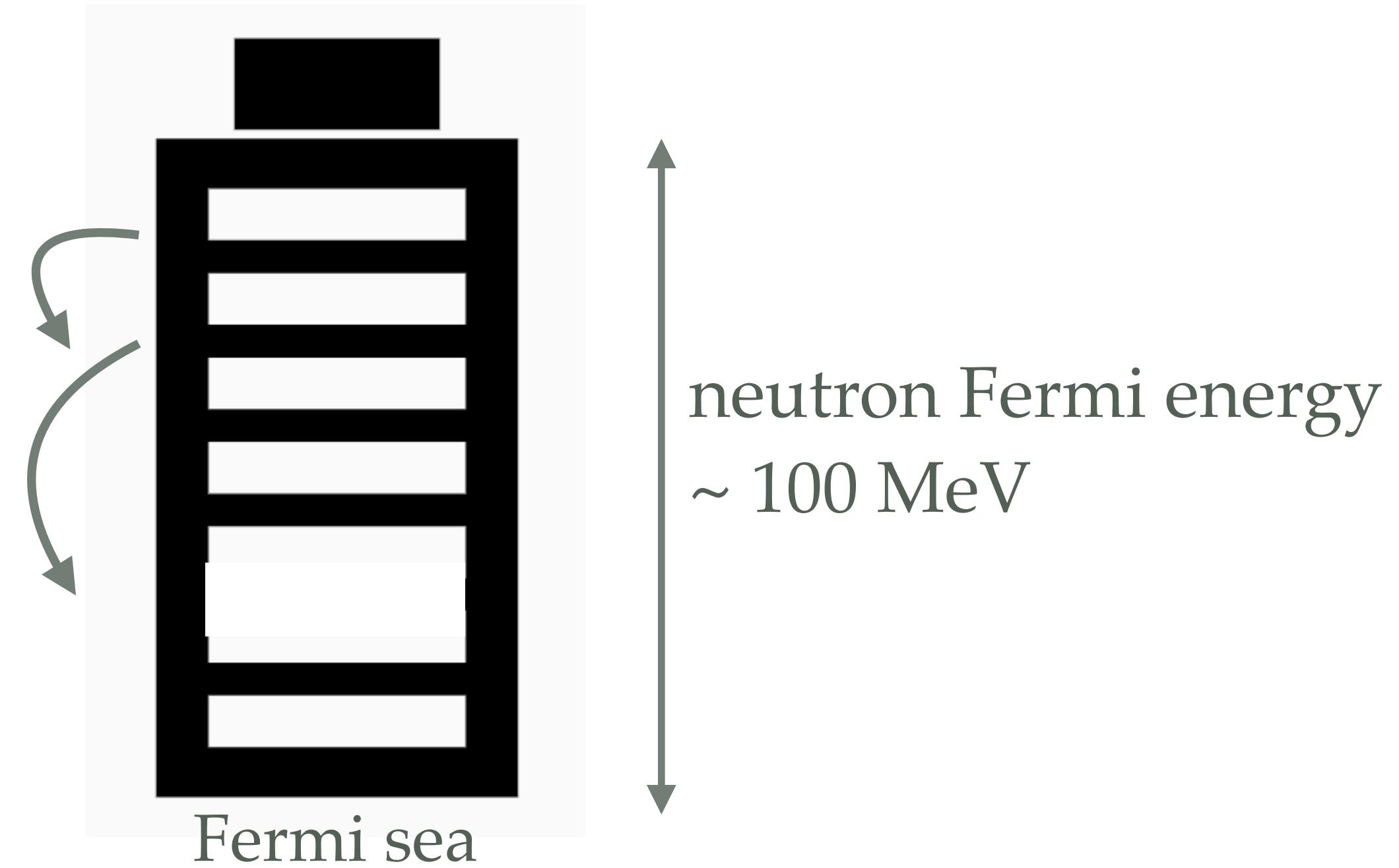
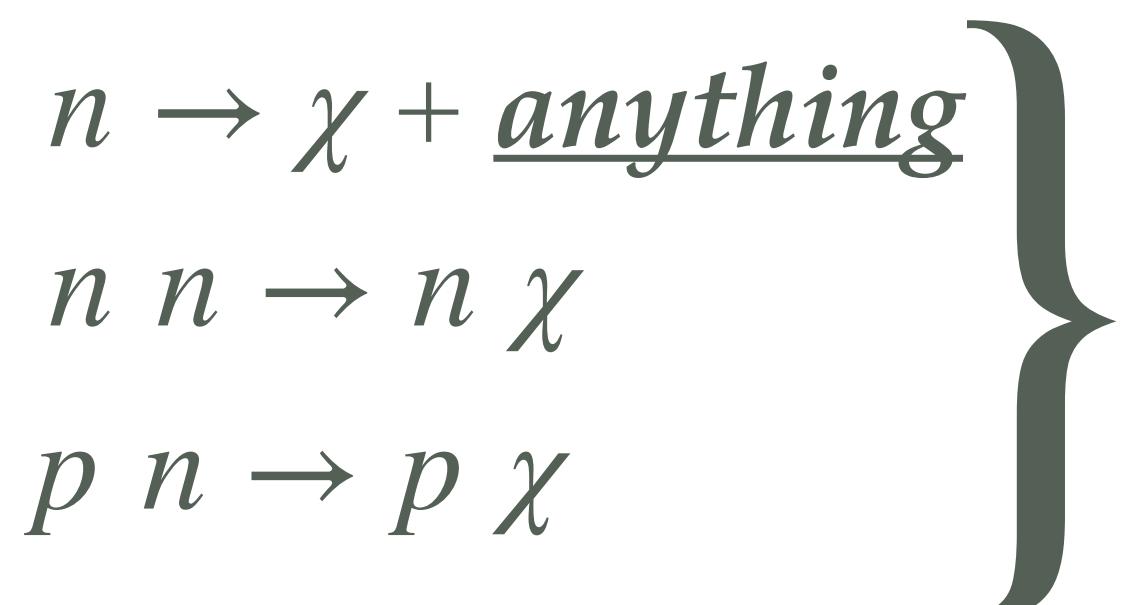
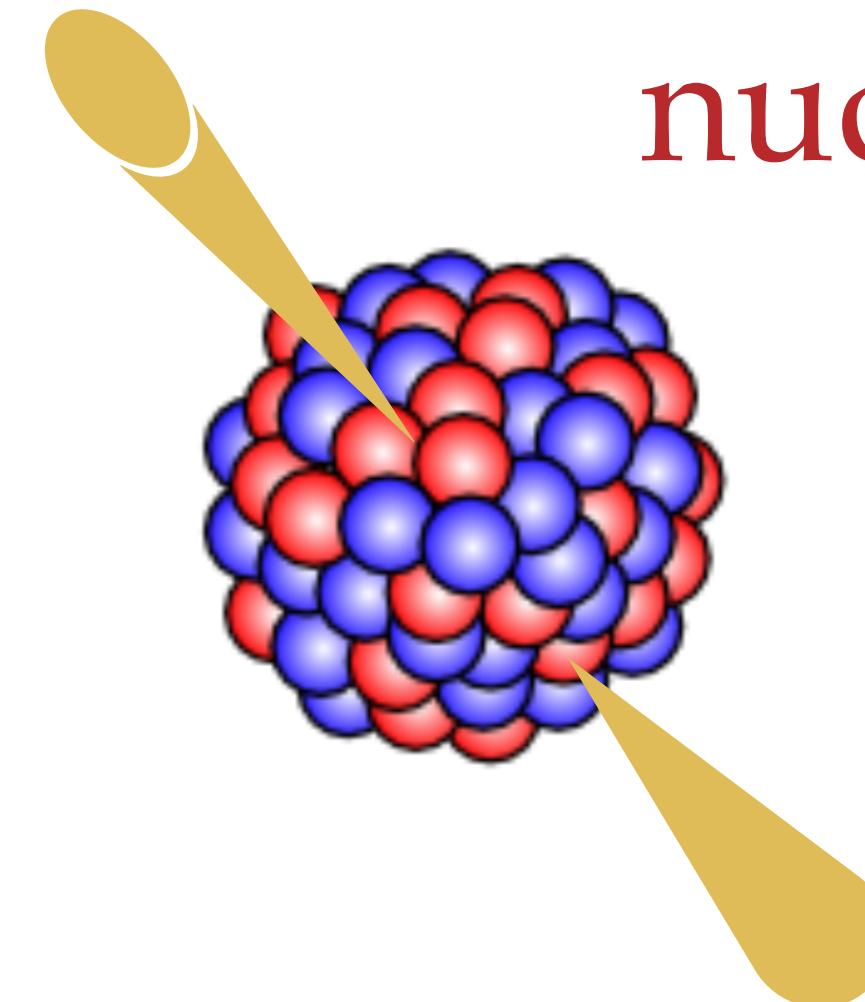
$$\Gamma_{\chi \rightarrow pe^- \bar{\nu}} = \frac{1}{9 \times 10^{22} \text{ s}} \left(\frac{\theta}{10^{-10}} \right)^2 \frac{F(Q_\chi/m_e)}{F(Q_n/m_e)}$$

$$\Gamma_{\chi \rightarrow n\gamma} \simeq \frac{1}{2200 \text{ s}} \left(\frac{\theta}{10^{-10}} \right)^2 \left| \frac{\Delta m}{10 \text{ MeV}} \right|^3$$

e or γ could “rewrite” reionization history by dumping EM energy in Dark Ages
(i.e. modify optical depth)

Probes: [3] neutron star temperatures

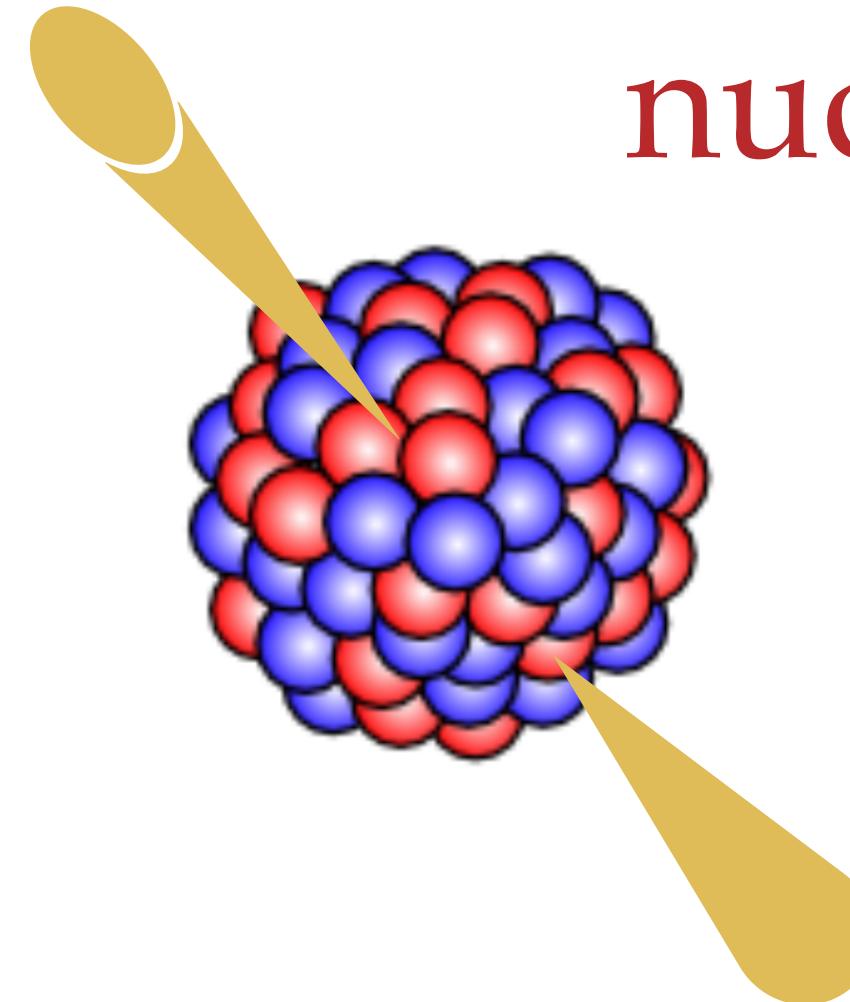
new heating mechanism:
nucleon “Auger effect”



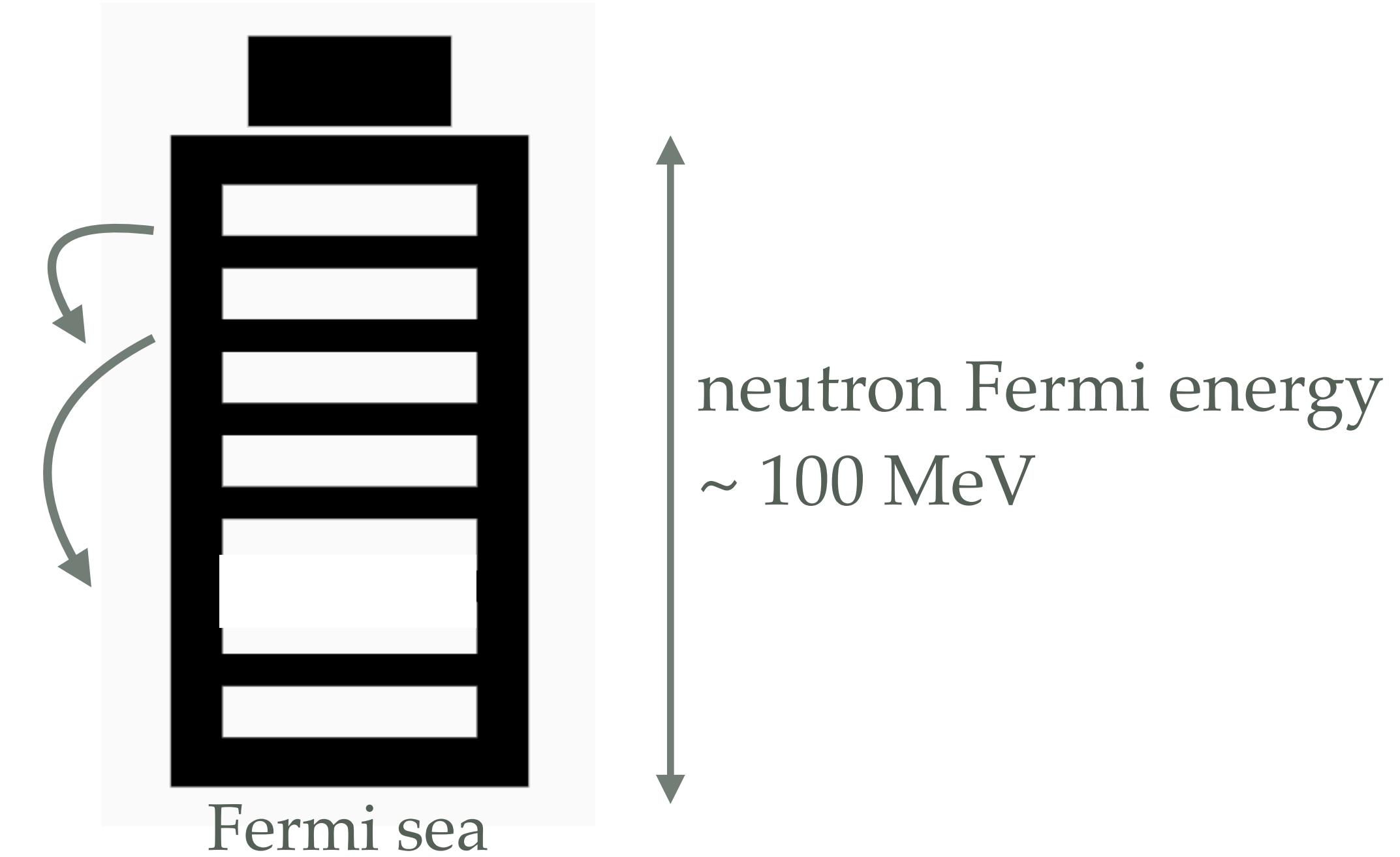
10^{57} neutrons
+
 10^{56} protons \Rightarrow explosive liberation of energy!

Probes: [3] neutron star temperatures

new heating mechanism:
nucleon “Auger effect”



$$\begin{aligned} n &\rightarrow \chi + \text{anything} \\ n \ n &\rightarrow n \ \chi \\ p \ n &\rightarrow p \ \chi \end{aligned} \quad \left. \right\}$$



*Hubble Space Telescope Nondetection of PSR J2144–3933: The Coldest Known Neutron Star**

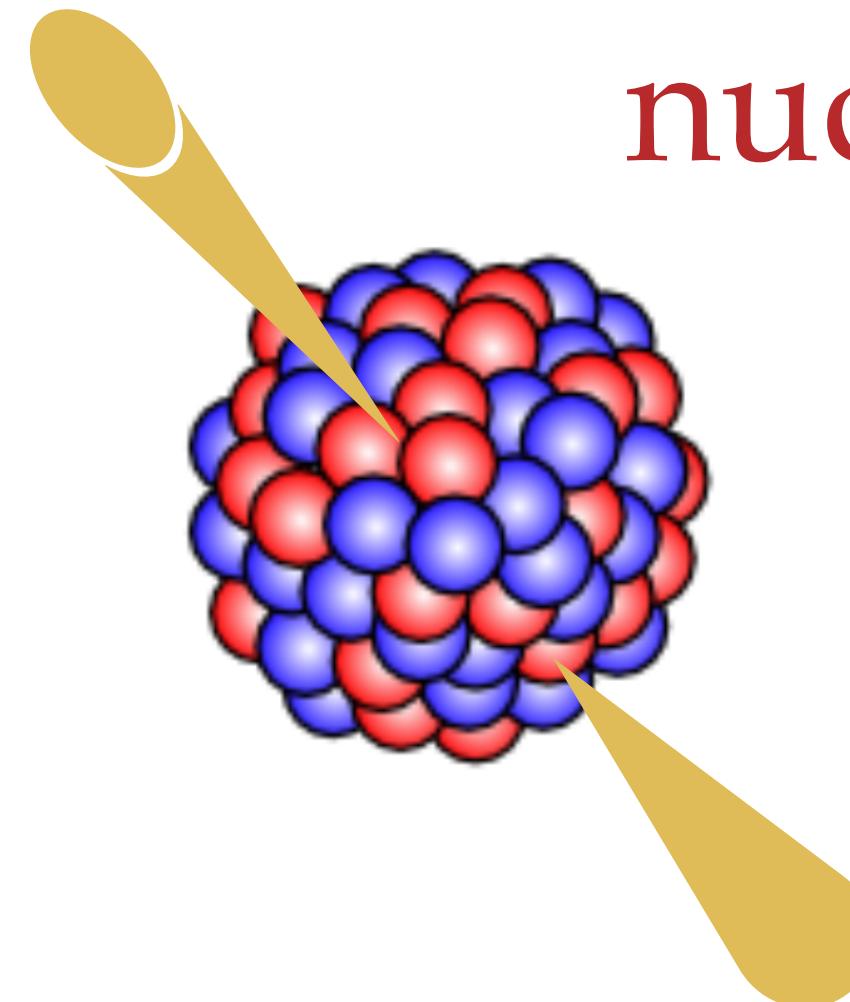
Sebastien Guillot^{1,2,3,8} , George G. Pavlov⁴ , Cristobal Reyes³ , Andreas Reisenegger³ , Luis E. Rodriguez⁵, Blagoy Rangelov⁶ , and Oleg Kargaltsev⁷

Suitable lab:

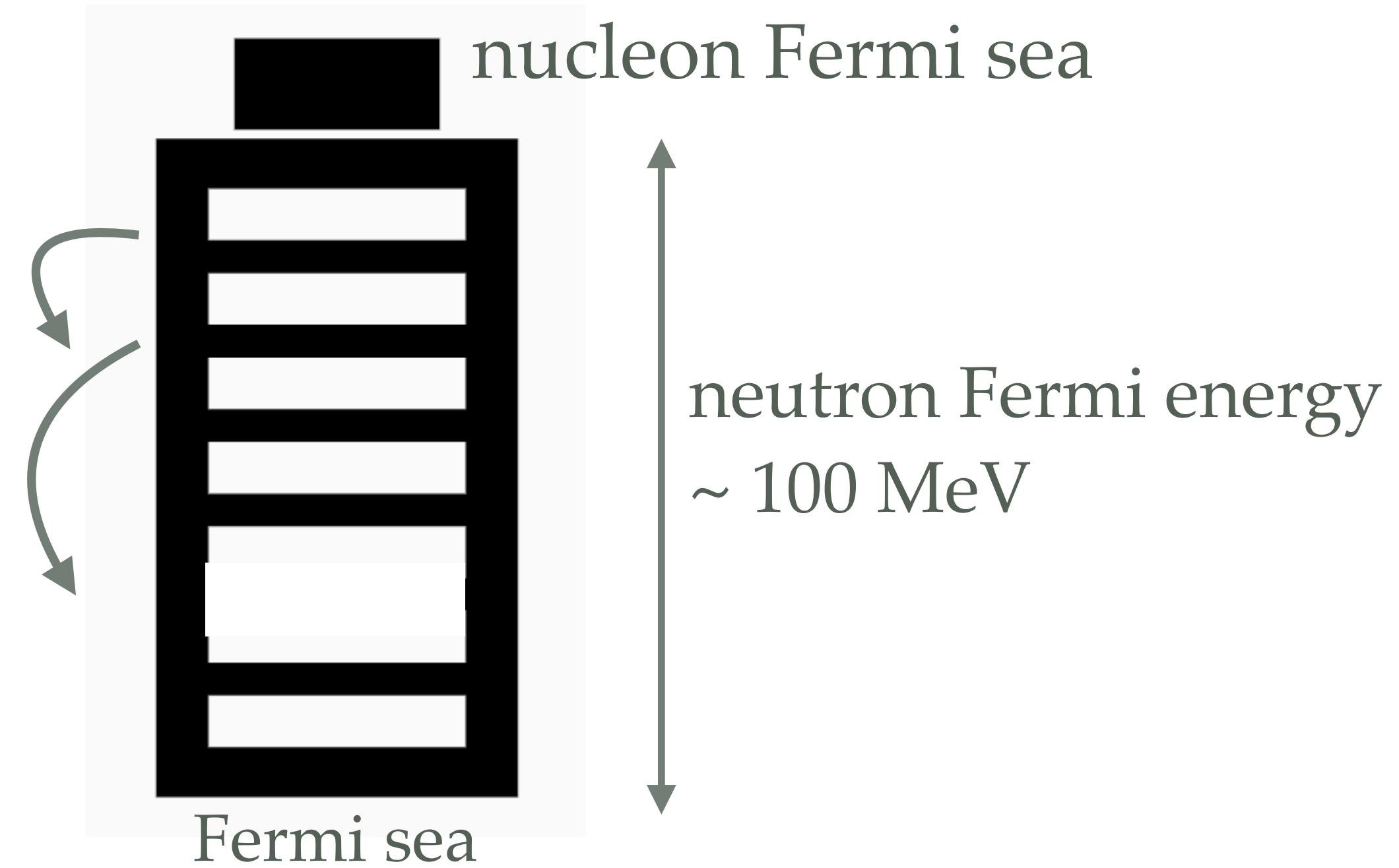
We report nondetections of the $\sim 3 \times 10^8$ yr old slow, isolated, rotation-powered pulsar PSR J2144–3933 in observations with the *Hubble Space Telescope* in one optical band (F475W) and two far-ultraviolet bands (F125LP and F140LP), yielding upper bounds $F_{\text{F475W}} < 22.7$ nJy, $F_{\text{F125LP}} < 5.9$ nJy, and $F_{\text{F140LP}} < 19.5$ nJy, at the pivot wavelengths 4940 Å, 1438 Å and 1528 Å, respectively. Assuming a blackbody spectrum, we deduce a conservative upper bound on the surface (unredshifted) temperature of the pulsar of $T < 42,000$ K. This makes

Probes: [3] neutron star temperatures

new heating mechanism:
nucleon “Auger effect”



$$\begin{aligned} n &\rightarrow \chi + \text{anything} \\ n \ n &\rightarrow n \ \chi \\ p \ n &\rightarrow p \ \chi \end{aligned} \quad \left. \right\}$$



Dark Kinetic Heating of Neutron Stars and an Infrared Window on WIMPs, SIMPs, and Pure Higgsinos

Masha Baryakhtar,¹ Joseph Bramante,¹ Shirley Weishi Li,² Tim Linden,² and Nirmal Raj³

¹*Perimeter Institute for Theoretical Physics, Waterloo, Ontario N2L 2Y5, Canada*

²*CCAPP and Department of Physics, The Ohio State University, Columbus, Ohio 43210, USA*

³*Department of Physics, University of Notre Dame, Notre Dame, Indiana 46556, USA*

(Received 10 April 2017; revised manuscript received 20 July 2017; published 26 September 2017)

We identify a largely model-independent signature of dark matter (DM) interactions with nucleons and electrons. DM in the local galactic halo, gravitationally accelerated to over half the speed of light, scatters against and deposits kinetic energy into neutron stars, heating them to infrared blackbody temperatures. The resulting radiation could potentially be detected by the James Webb Space Telescope, the Thirty Meter Telescope, or the European Extremely Large Telescope. This mechanism also produces optical emission

Future lab:

optimized for
~2000 K

Aside: more fundamental physics with UV, optical, IR telescope missions

31801 (2017)

PHYSICAL REVIEW LETTERS

week
29 SEPT

Hubble Space Telescope Nondetection of PSR J2144–3933: The Coldest Known Neutron Star*

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Scattering searches for dark matter in subhalos: neutron stars, cosmic rays, and old rocks

Joseph Bramante,^{1, 2,*} Bradley J. Kavanagh,^{3, †} and Nirmal Raj^{4, ‡}

submitted to PRL,
arXiv: 2109.04582

see also:

N. Raj, P. Tanedo, H-B. Yu PRD 2017

EFT

J. Acevedo, J. Bramante, R. Leane, N. Raj, JCAP 2020 NS crust/pasta

A. Joglekar, N. Raj, P. Tanedo, H-B. Yu PLB 2020 NS electrons

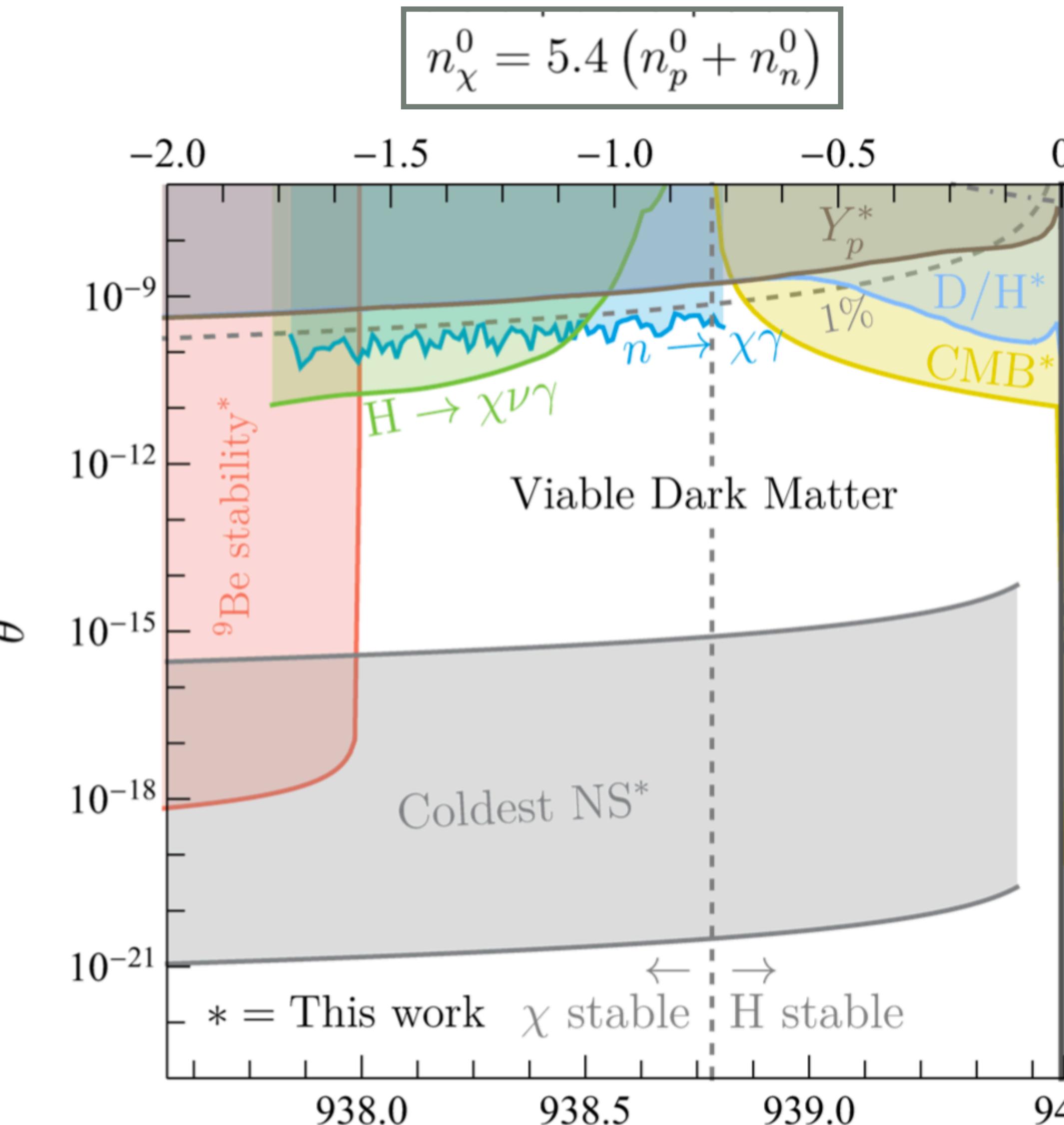
A. Joglekar, N. Raj, P. Tanedo, H-B. Yu PRD 2020 NS electrons EFT

R. Garani, A. Gupta, N. Raj, PRD 2021 DM refrigeration

Constraints

$n \rightarrow \chi \gamma$
open

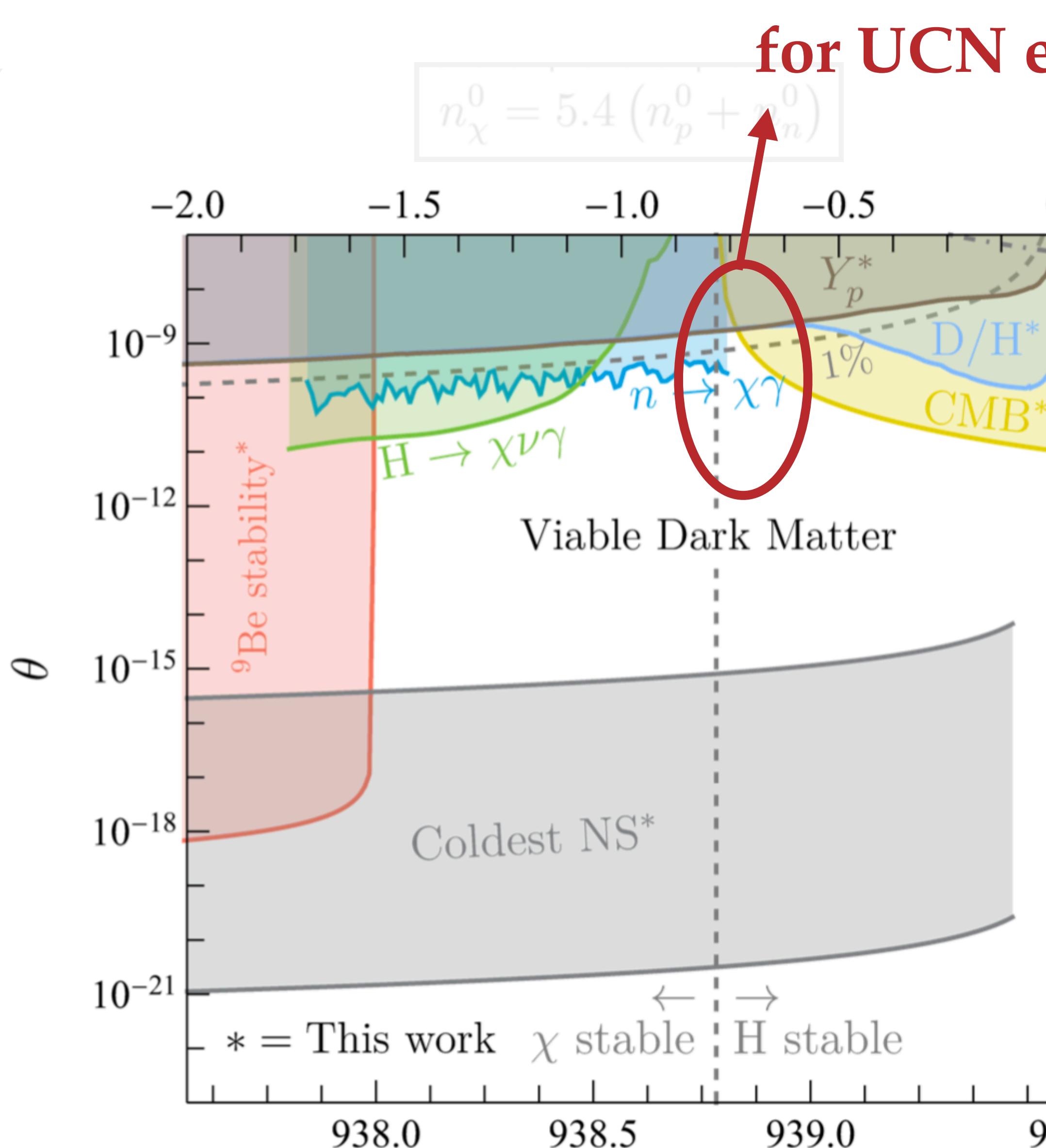
longer
life



- BBN data: $Y_p = 0.245 \pm 0.004$,
 $D/H = (2.55 \pm 0.03) \times 10^{-5}$,
 ${}^3\text{He}/H = (1.0 \pm 0.5) \times 10^{-5}$,
- CMB limit: $f_\chi/\tau_\chi \lesssim 10^{-25} \text{ s}^{-1}$
T. R. Slatyer, Physical Review D **87** (2013), [10.1103/physrevd.87.123513](https://doi.org/10.1103/physrevd.87.123513).
J. M. Cline and P. Scott, JCAP **03**, 044 (2013), [Erratum: JCAP 05, E01 (2013)], arXiv:1301.5908 [astro-ph.CO].
- $n \rightarrow \chi \gamma$ direct search: 1802.01595 [nucl-ex]
- $H \rightarrow \chi \nu \gamma$: Borexino recast by McKeen, Pospelov (2003.02270)
- ${}^9\text{Be} \rightarrow 2 {}^4\text{He} + \chi$:
Limited by: $\tau_{{}^9\text{Be}} \sim 4 \times 10^{10} \text{ yr} \left(\frac{10^{-19}}{\theta}\right)^2 \left(\frac{1 \text{ MeV}}{Q_{{}^9\text{Be}}}\right)^{3/2}$
 $< 3 \times 10^9 \text{ yr}$ in metal-poor stars
- NS: J2144-3933

Constraints

$n \rightarrow \chi\gamma$
open



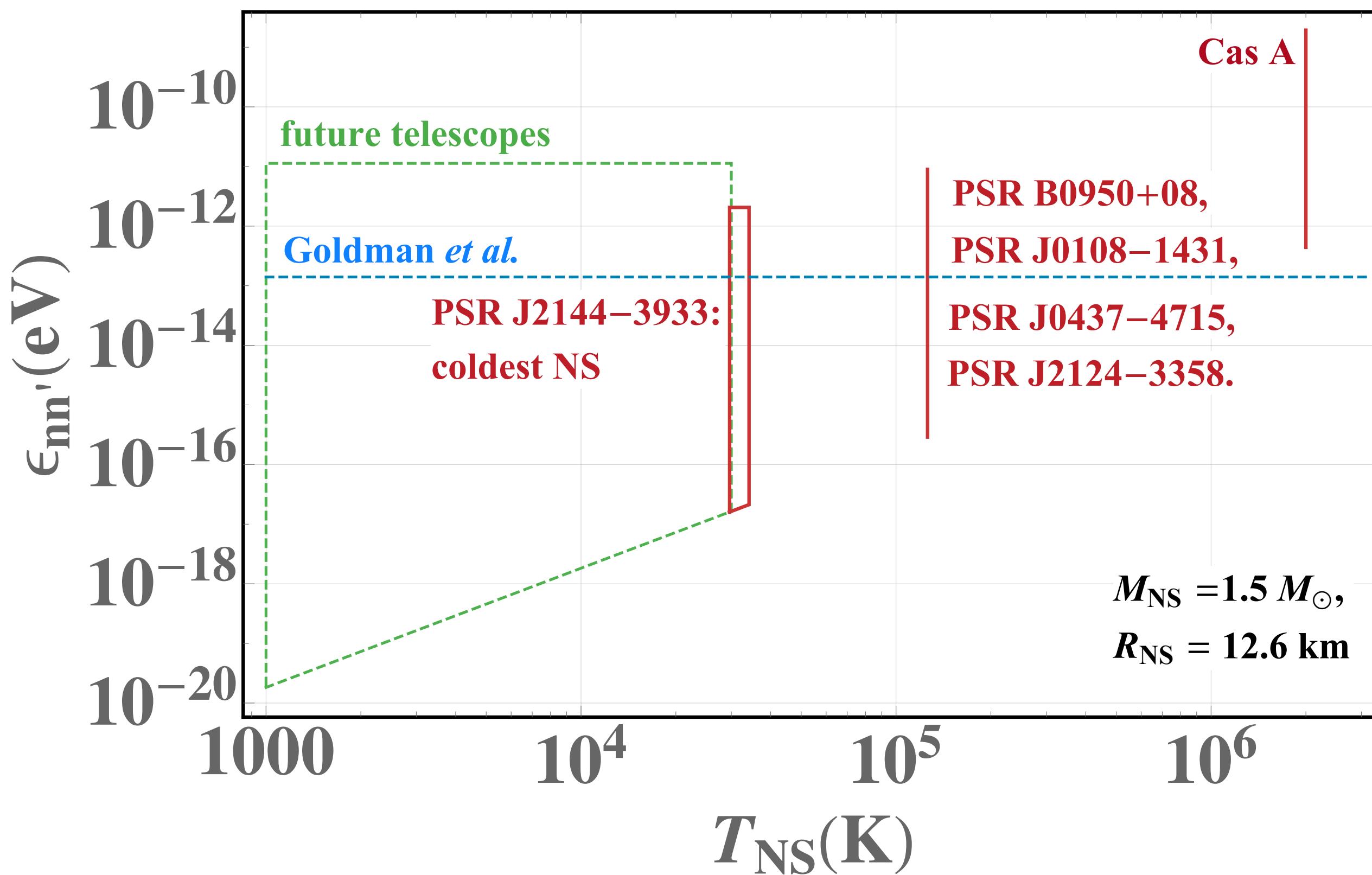
100 keV “neutron lifetime puzzle” window, for UCN experimentalists to target!

- CMB limit: $f_\chi/\tau_\chi \lesssim 10^{-25} \text{ s}^{-1}$
T. R. Slatyer, Physical Review D 87 (2013), 10.1103/physrevd.87.123513.
J. M. Cline and P. Scott, JCAP 03, 044 (2013), [Erratum: JCAP 05, E01 (2013)], arXiv:1301.5908 [astro-ph.CO].
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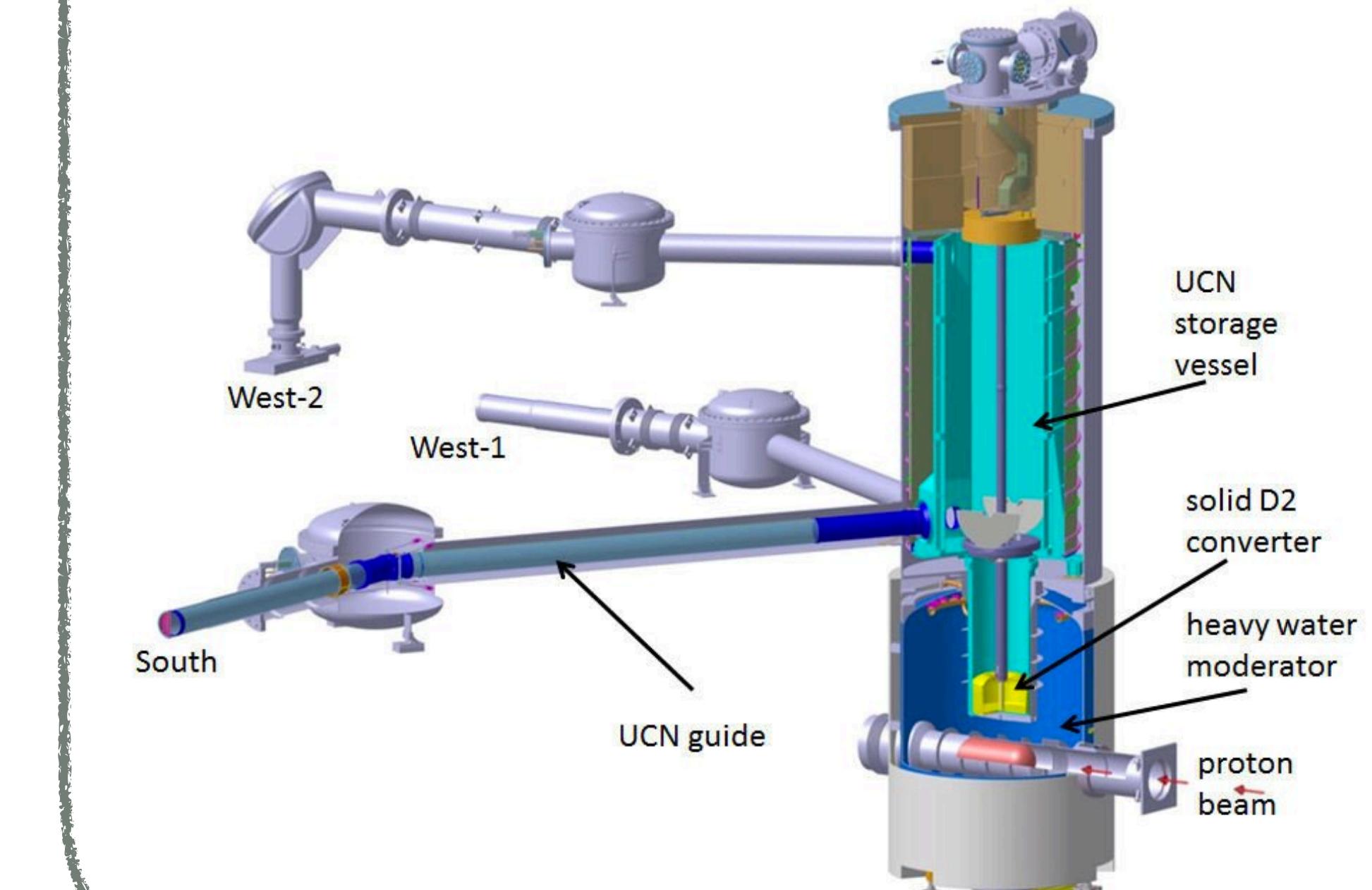
Constraints: NS heating

NS energy per baryon

neutron star heating: $|m_n - m_{n'}| \lesssim \mathcal{O}(10 \text{ MeV})$



What about terrestrial
ultra-cold neutron searches?



ceilings: neutron conversions stop within NS lifetime

NB. neutron lifetime anomaly explained by $\epsilon_{nn'} \sim 10^{-8} \text{ eV}$ (Berezhiani 2018)

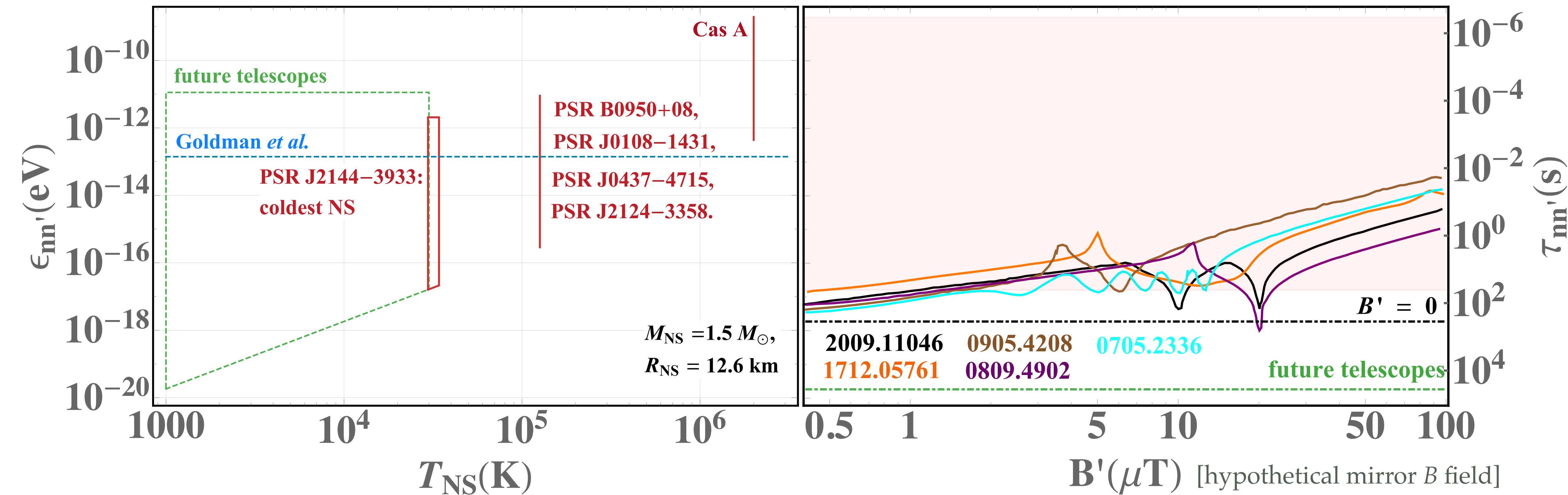
Constraints: NS heating

NS energy per baryon

Zeeman from Earth's B field

neutron star heating: $|m_n - m_{n'}| \lesssim \mathcal{O}(10 \text{ MeV})$

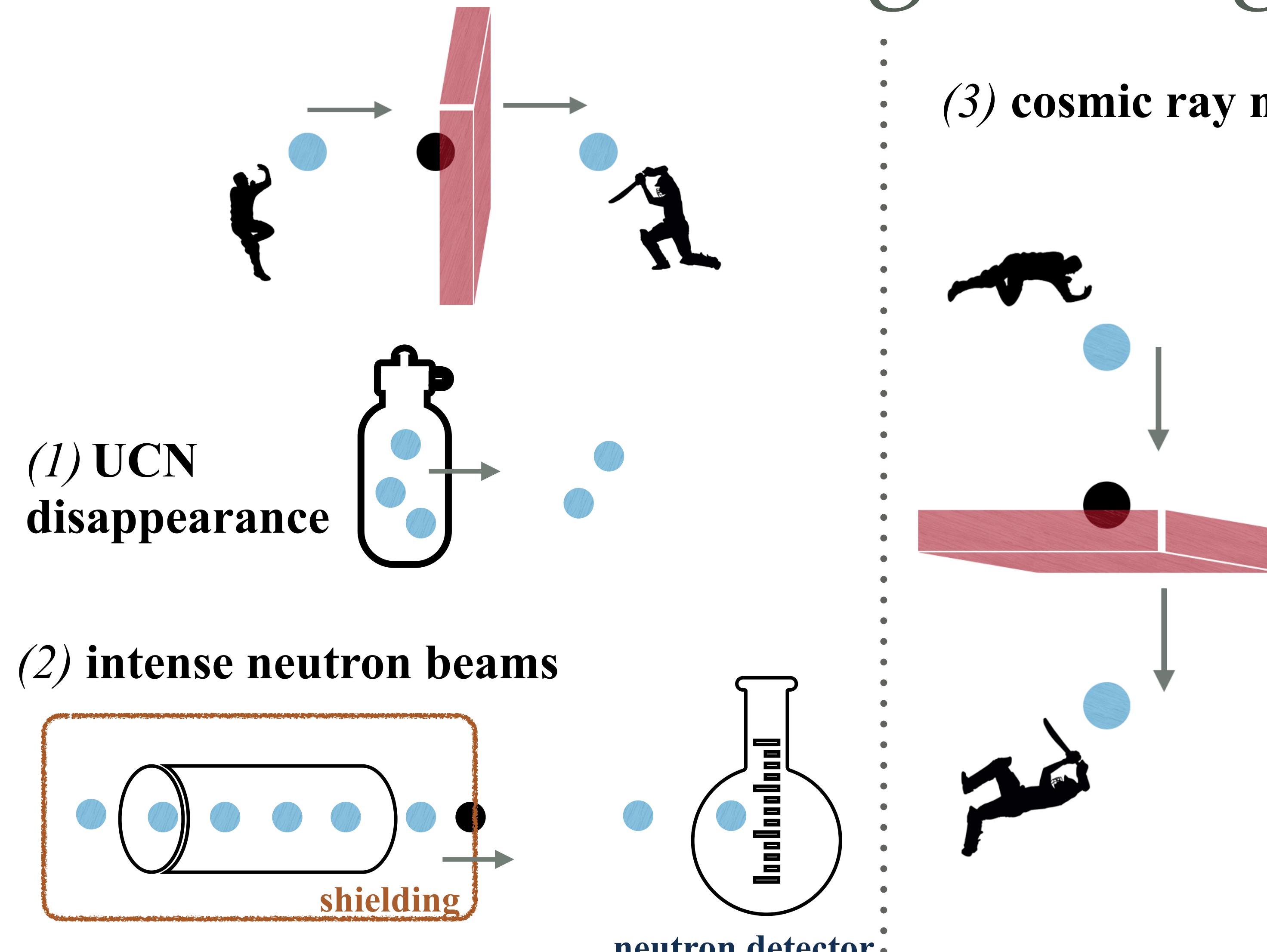
UCN searches: $|m_n - m_{n'}| < 10^{-18} \text{ MeV}$



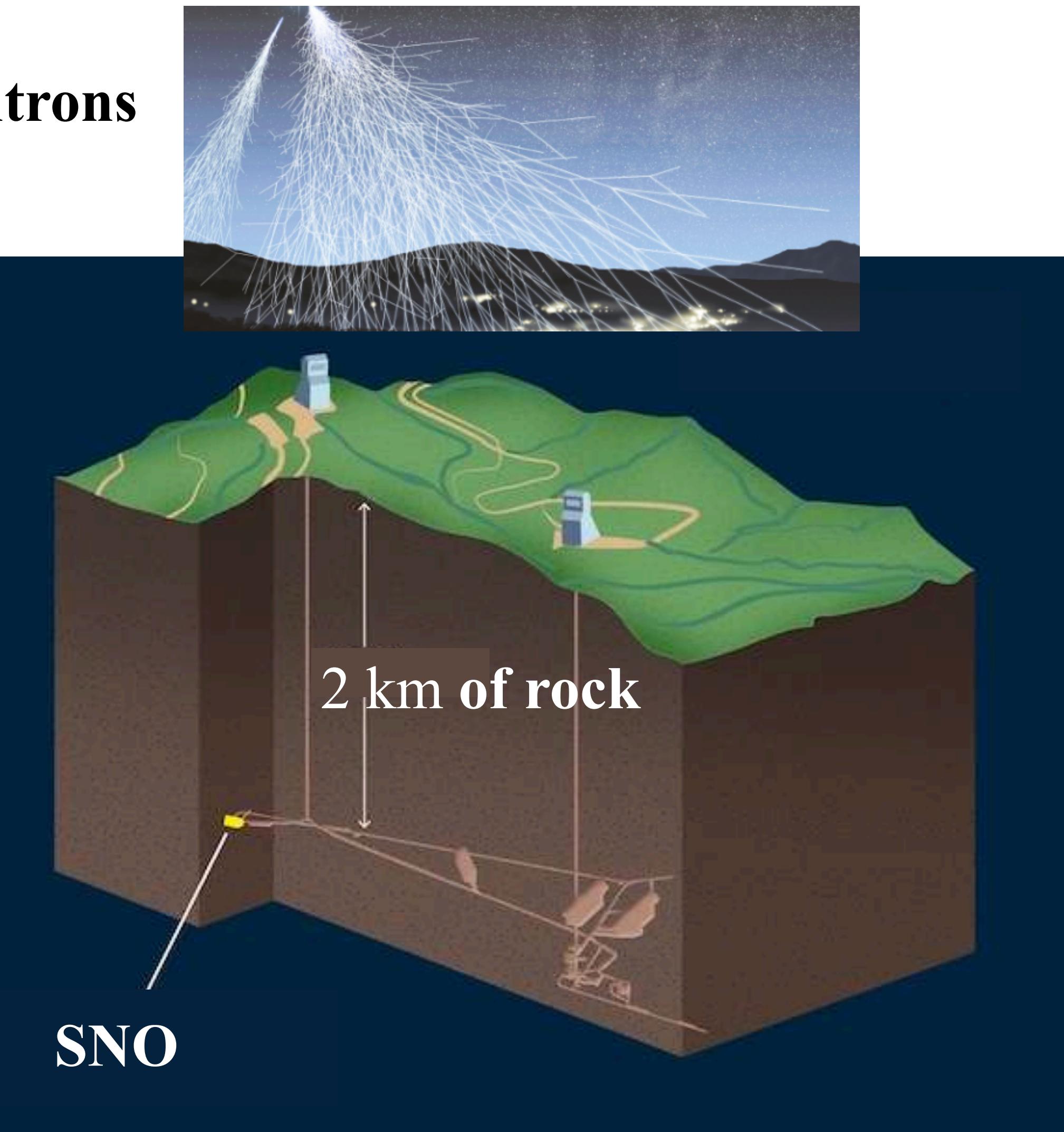
ceilings: neutron conversions stop within NS lifetime

NB. neutron lifetime anomaly explained by $\epsilon_{nn'} \sim 10^{-8} \text{ eV}$ (*Berezhiani 2018*)

Neutrons shining through a wall



e.g., IsoDAR at Yemilab:
Conrad, *et al.*, 2110.10635



ongoing work with
M. Hostert, D. McKeen, M. Pospelov

Highlights

- Cosmology (BBN + CMB) stringently limits dark neutron explanation of neutron lifetime puzzle.
 - small 100 keV-ish window left for UCN experiments to target!
- very slow exotic neutron decays => explosive heating of neutron stars.
- Heavier-than-neutron dark neutrons (see back-up slides): cosmology sole probe.

Thank you! Questions?

Back-up slides

Constraining neutron conversions

heating rate
cooling rate
(blackbody emission)

$$\int_{\text{NS}} d^3r \, n_n(\mathbf{r}) \dot{E}_{n'}(\mathbf{r}) < 4\pi R_{\text{NS}}^2 \sigma_{\text{SB}} T_{\text{NS}}^4$$

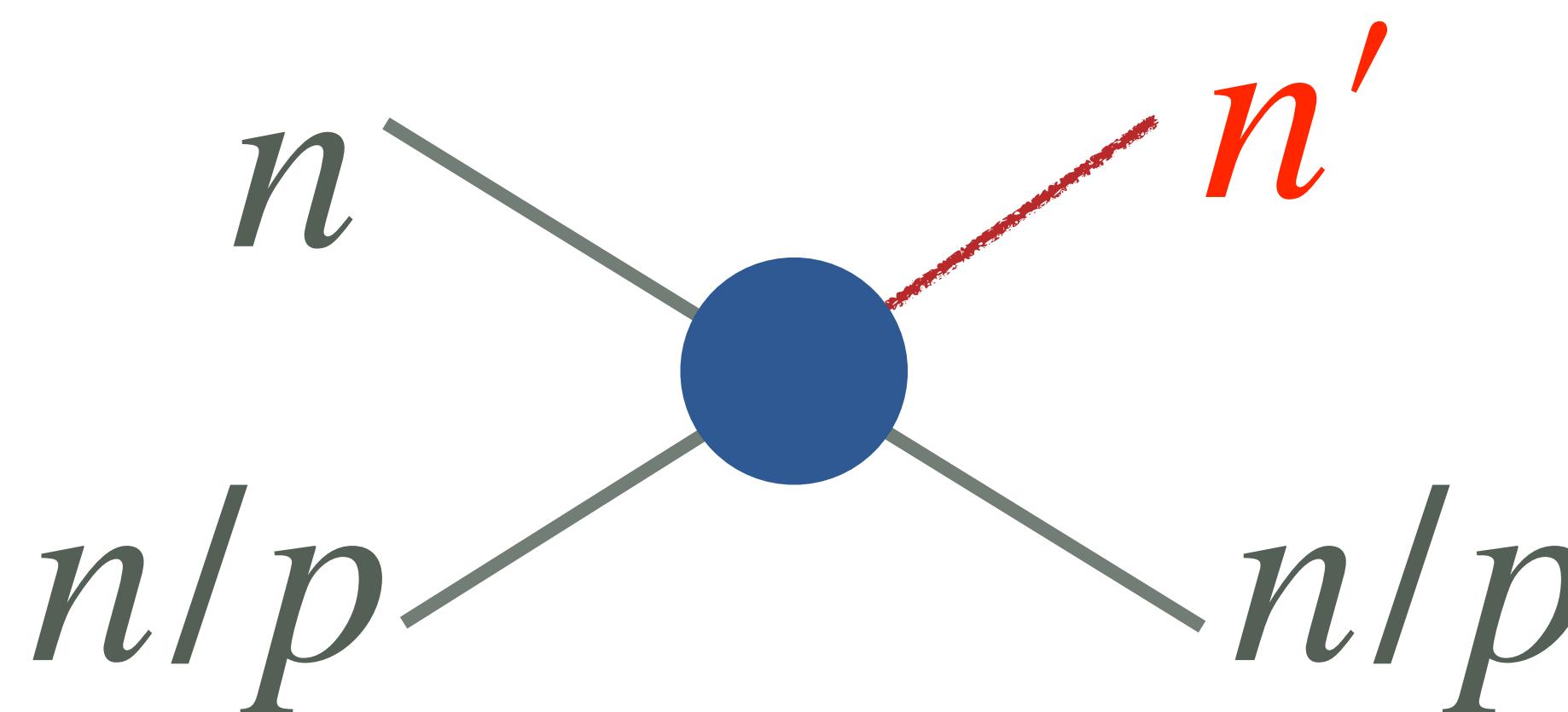
neutron number density
energy release rate

Conversions to dark neutrons

$$H = \begin{pmatrix} m_n + \Delta E & \epsilon_{nn'} \\ \epsilon_{nn'} & m_{n'} \end{pmatrix}$$

medium-dependent splitting
e.g. neutron star nuclear self-energies, 10—100 MeV

$$\sigma_{n'N} \simeq g_N \left(\frac{\epsilon_{nn'}}{\Delta E} \right)^2 \sigma_{nN \rightarrow nN}$$



$$\sigma_{nn \rightarrow nn} \simeq \frac{1}{4} \times \frac{16\pi}{m_N^2 v^2} \sin^2 \delta_S,$$
$$\sigma_{np \rightarrow np} \simeq \frac{1}{4} \times \frac{16\pi}{m_N^2 v^2} (\sin^2 \delta_S + 3 \sin^2 \delta_T)$$

energy-dependent
phase shifts
from nuclear potential models
(<https://nn-online.org/>)

Conversions to dark neutrons

$$\dot{E}_{n'} = \sum_{N=n,p} f_N n_N \left\langle \left(\tilde{\mu}_n - \frac{p_{n'}^2}{2m_{n'}} \right) \sigma_{n'N} v \right\rangle_{p_N > p_{F_N}}$$

symmetry factor

energy release rate

number density*

neutron chemical potential*

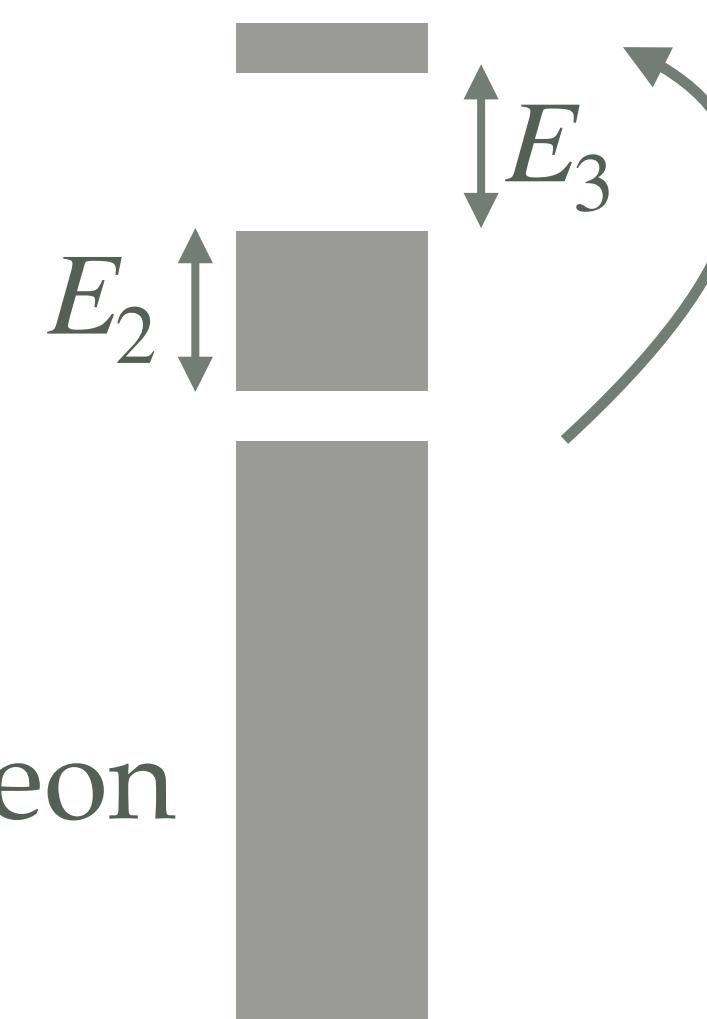
Pauli blocking condition

3 sources of energy:

scattering neutron
Fermi sea



spectator nucleon
Fermi sea



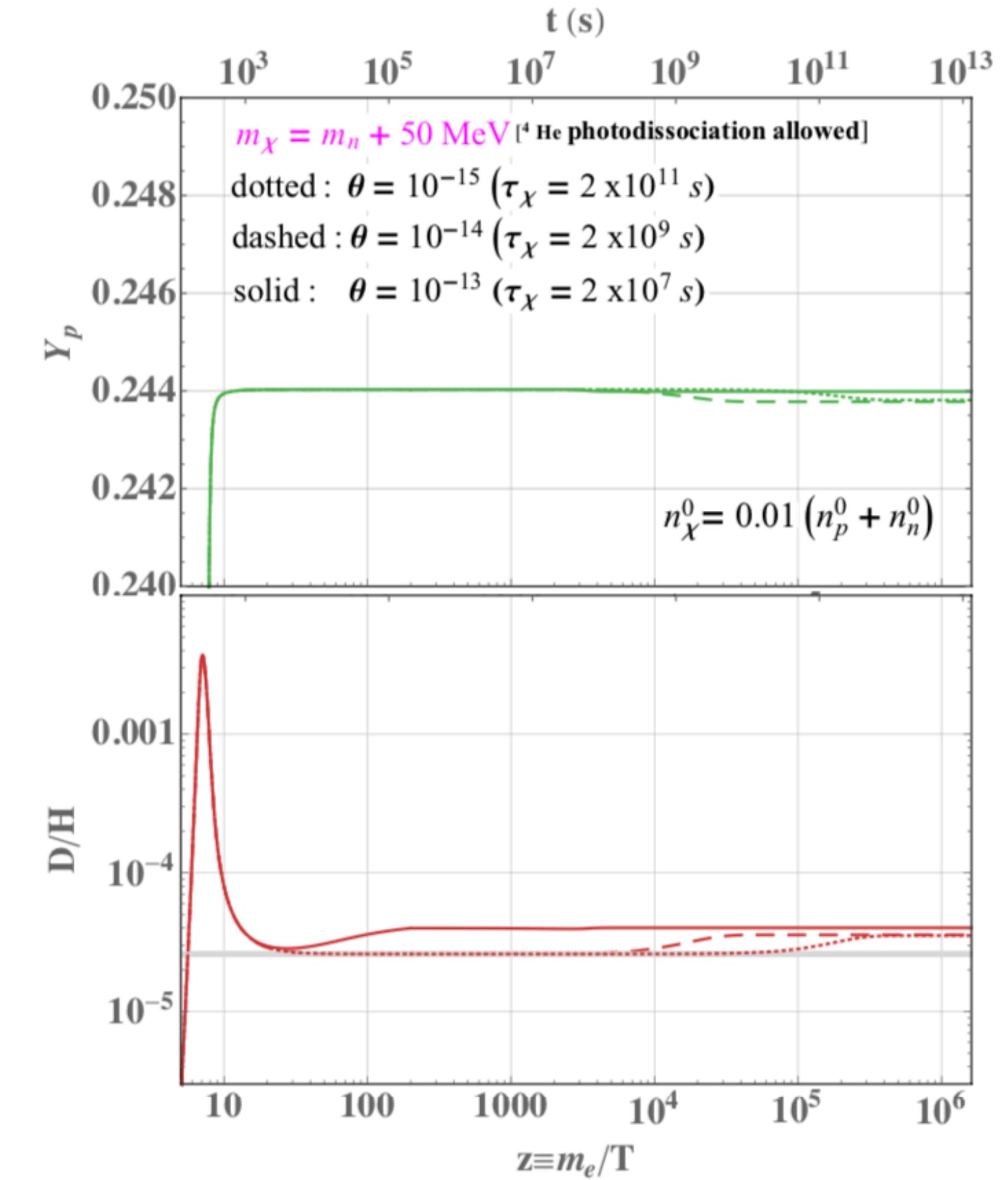
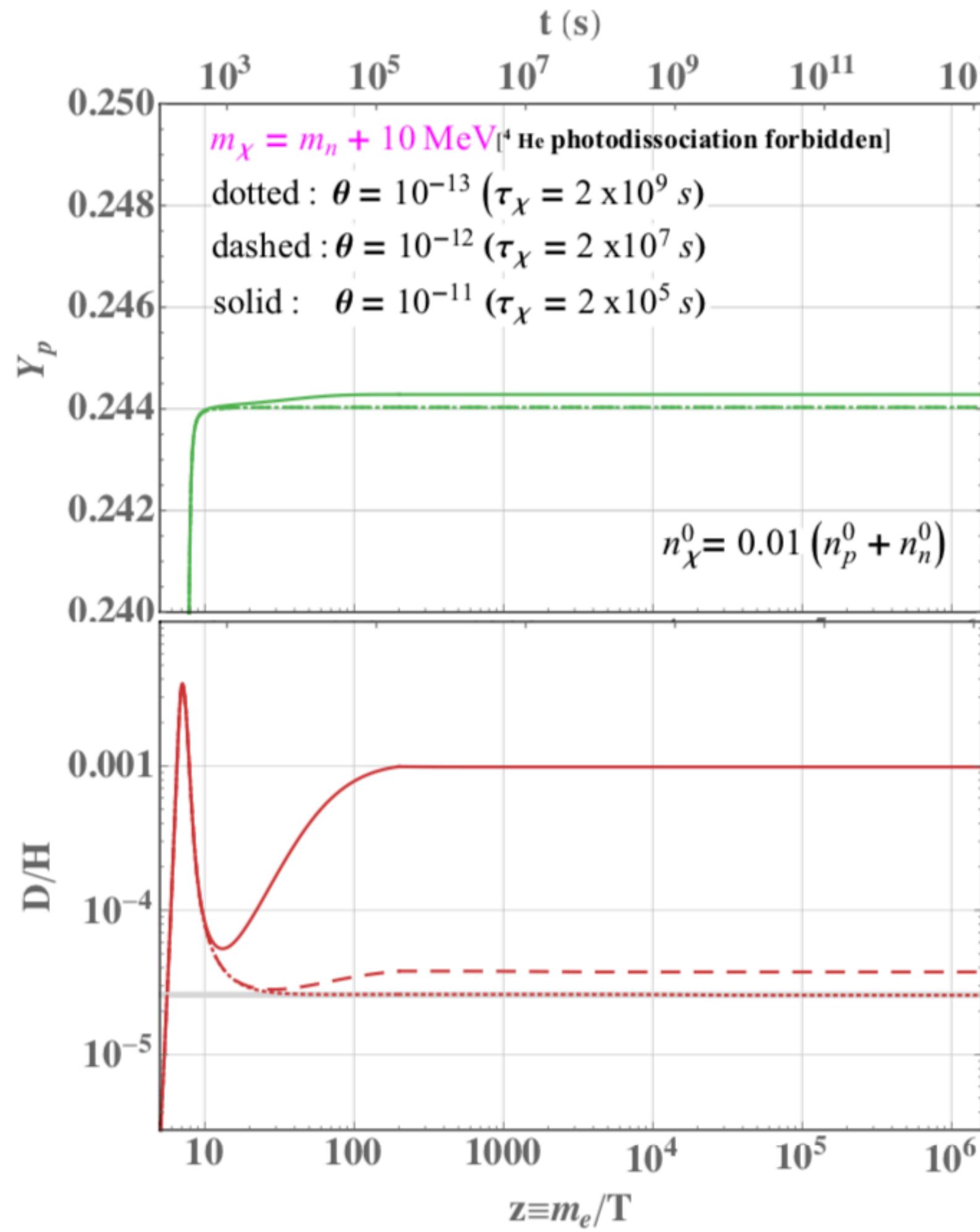
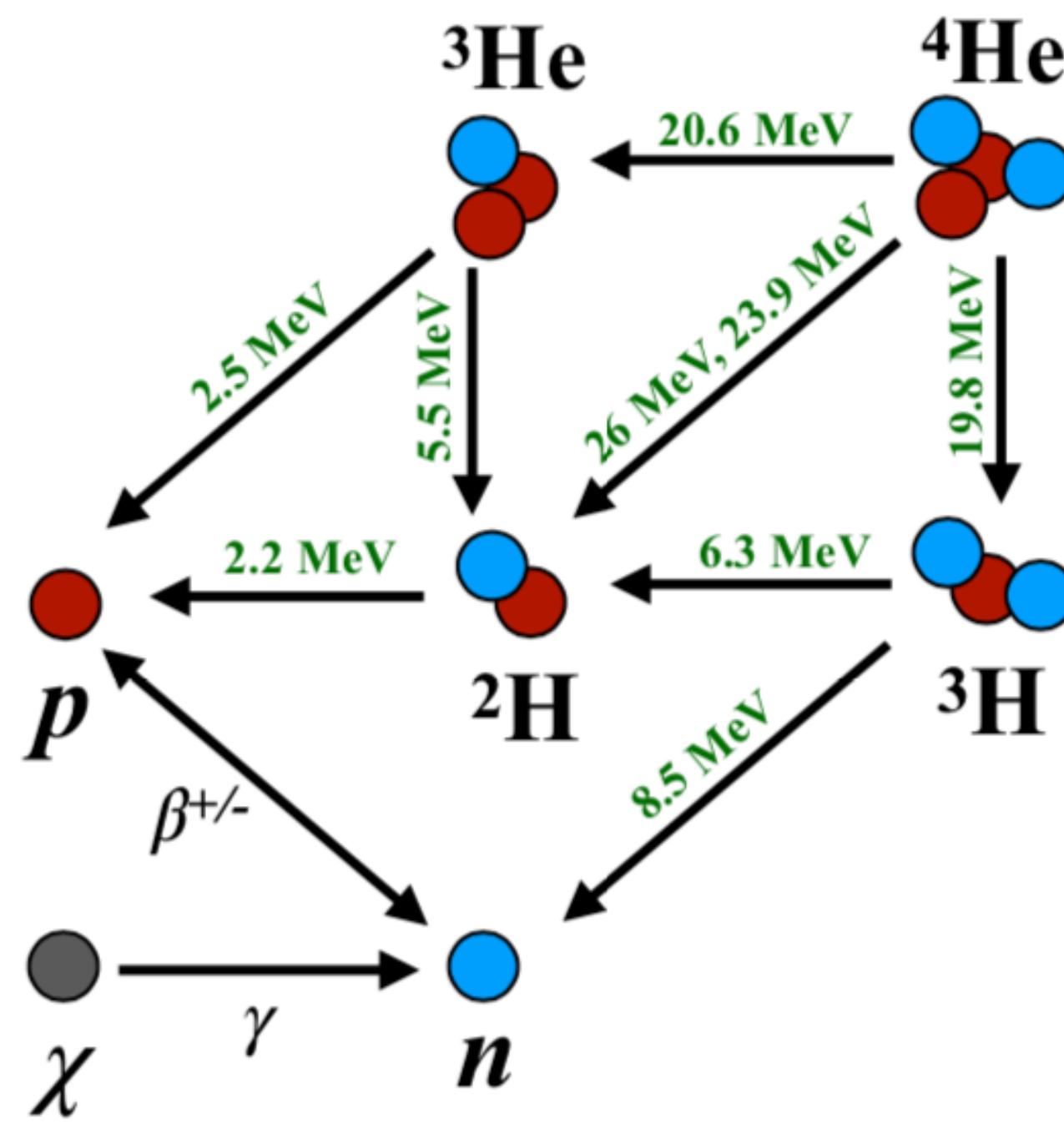
Amusement

proton spectators
(~ 10% of NS nucleons)
supply more heat!

less Pauli-blocked,
greater cross section

* determined from high-density equation of state + NS mass & radius,
in practice used Brussels-Montreal BSk24 with $M_{\text{NS}} = 1.5 M_{\odot}$, $R_{\text{NS}} = 12.6$ km

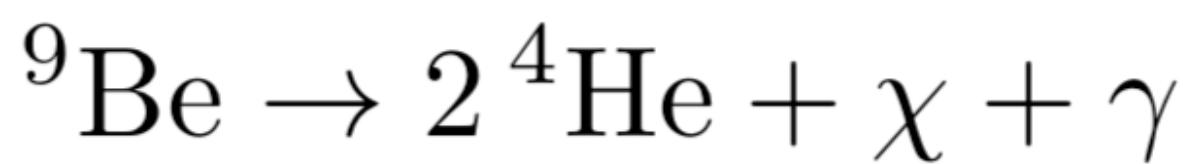
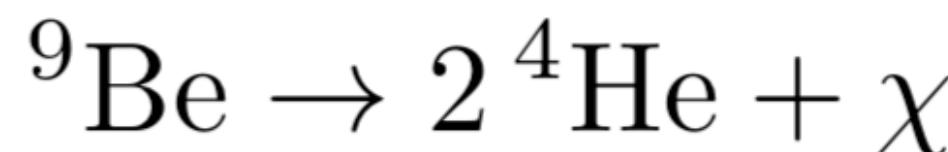
Signals: photodissociation post-nucleosynthesis



Probes

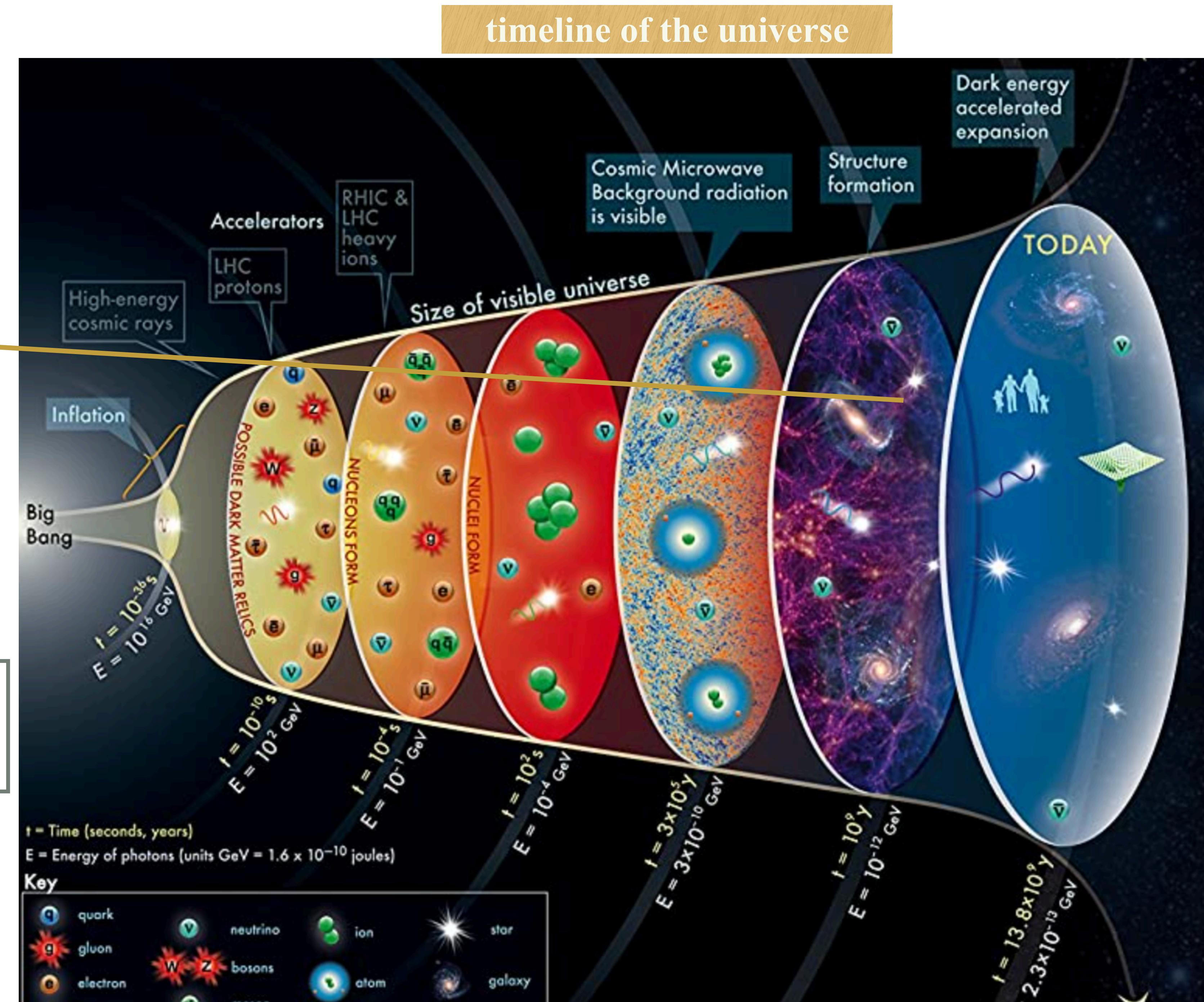
(4) ancient metal-poor stars
(~3 Gyr old; ${}^9\text{Be}$ observed)

Via



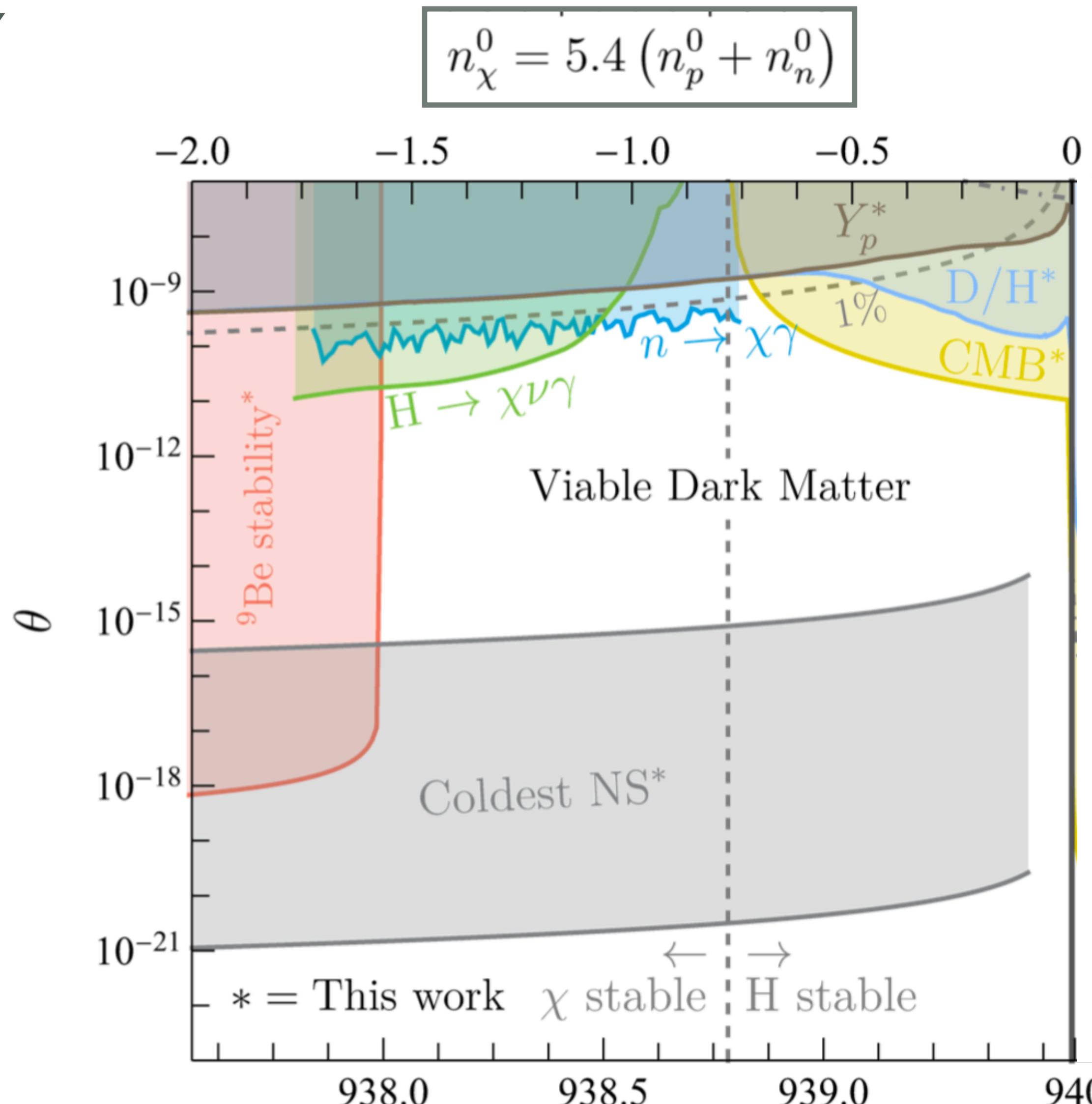
beryllium-9 lifetime:

$$4 \times 10^{10} \text{ yr} \left(\frac{10^{-19}}{\theta} \right)^2 \left(\frac{1 \text{ MeV}}{Q_{{}^9\text{Be}}} \right)^{3/2}$$



Constraints

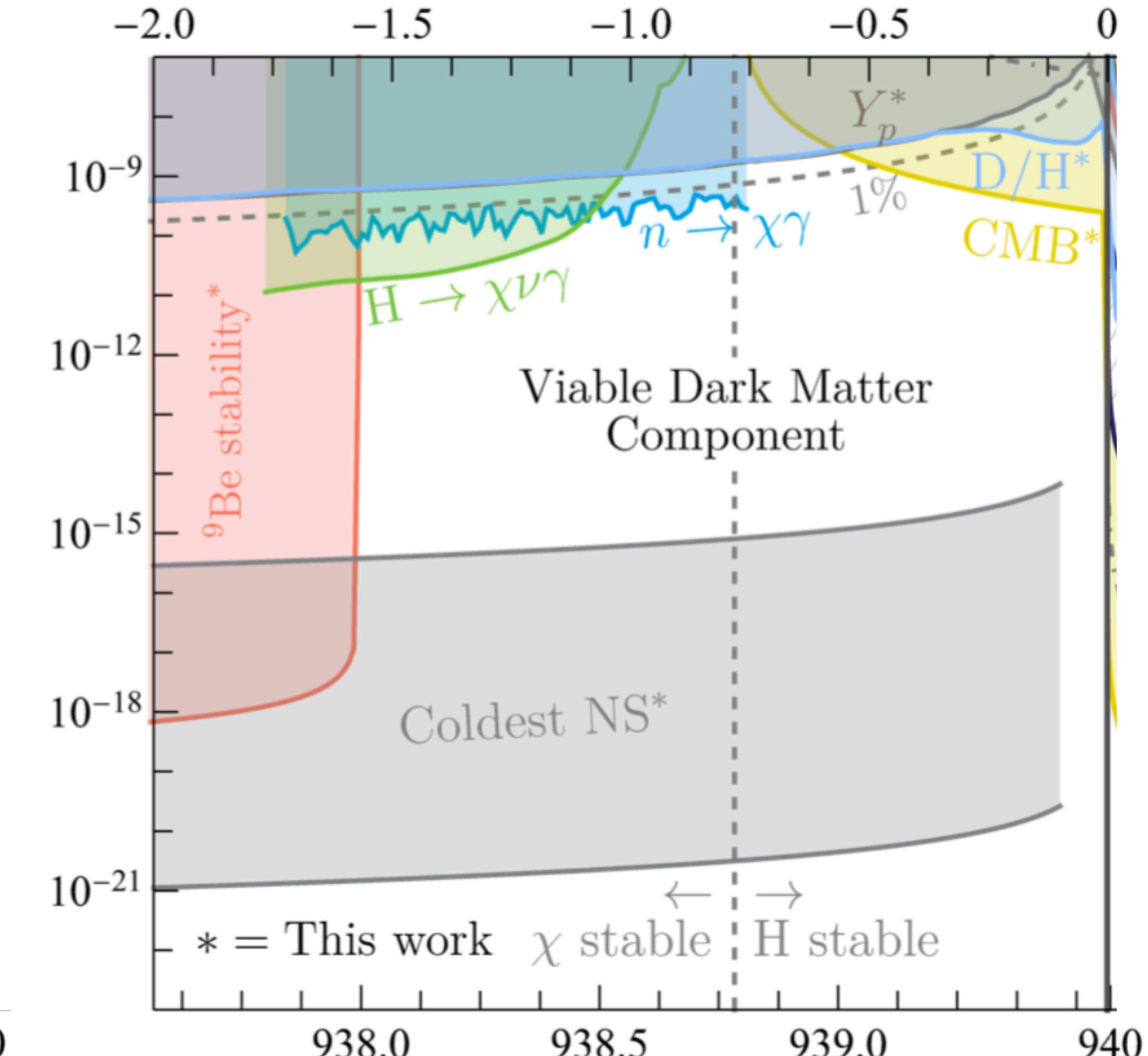
$n \rightarrow \chi\gamma$
open



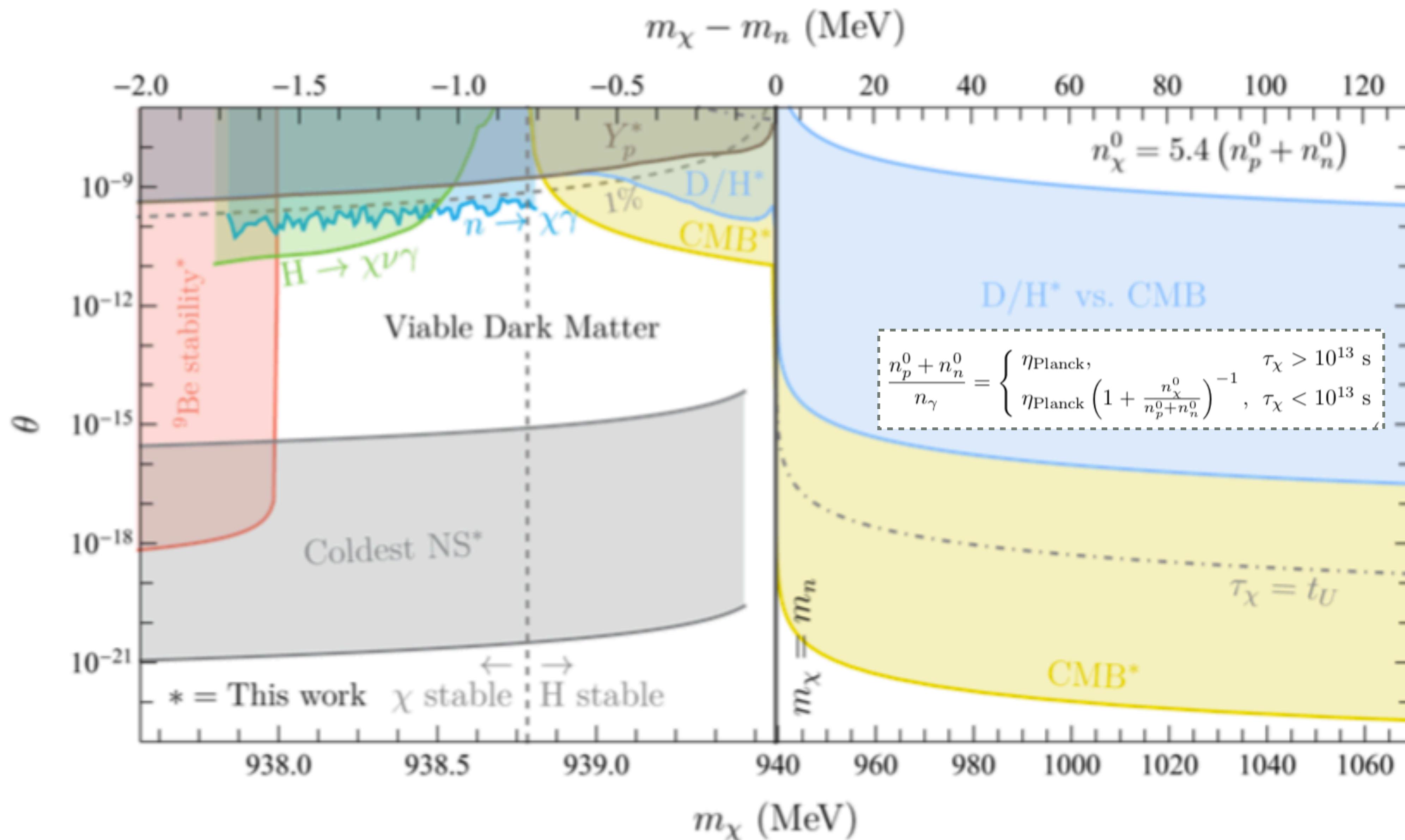
$$n_\chi^0 = 5.4 (n_p^0 + n_n^0)$$

$$n_\chi^0 = 0.01 (n_p^0 + n_n^0)$$

longer
life



Constraints: χ all the dark matter



Constraints: χ percent-level dark matter

